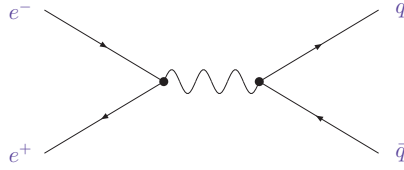


Tutorial II

IMSc Spring School on High Energy Physics 2025

The amplitude for electron-positron annihilation and subsequent quark-anti-quark production is given to leading order by the following diagram:



The photon has momentum q , $q^2 = Q^2 > 0$ and the quark and anti-quark carry the momenta p , $p^2 = 0$ and p_1 , $p_1^2 = 0$.

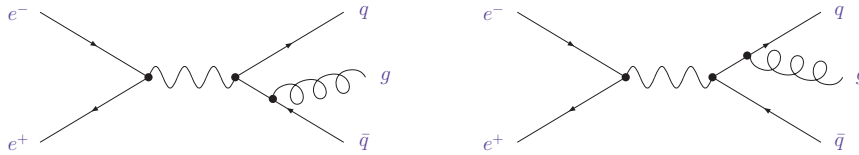
- Calculate the total cross section for this reaction.

$$e^- + e^+ \rightarrow \gamma(q) \rightarrow q(p) + \bar{q}(p_1). \quad (1)$$

Next, we want to consider the radiative corrections in Quantum Chromodynamics to order α_s . The amplitude for electron-positron annihilation and production of a quark-anti-quark pair and additional radiation of a gluon with momentum k ,

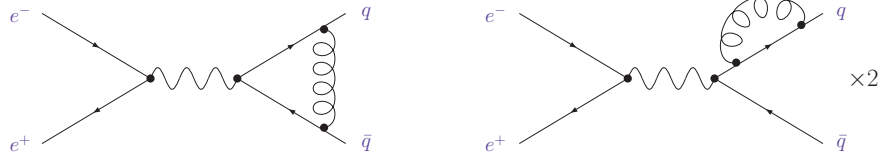
$$e^- + e^+ \rightarrow \gamma(q) \rightarrow q(p) + \bar{q}(p_1) + g(p_2), \quad (2)$$

is given by the following diagrams:



- This process can be described with the help of the scaling variable $x = (2p \cdot q)/Q^2$. What is the physical interpretation of x ?
- Calculate the cross section $d\sigma_T$ for this reaction for transversely polarized photons as a function of the scaling variable $x = (2p \cdot q)/Q^2$ and the photon momentum $q^2 = Q^2$.
- Why is the result divergent in the infrared?

- What is the role of the following set of diagrams?



Hint: The squared amplitude $\mathcal{A}$ for the reaction $\gamma(q, \mu) \rightarrow q(p) + \bar{q}(p_1) + g(p_2, \mu_1)$ is given by

$$\sum_{spins} |\mathcal{A}|^2 = (e^2 g_s^2 C_F) g^{\mu\nu} g^{\mu_1\mu_2} \quad (3)$$

$$\left\{ \text{tr}(\not{p}_1 \gamma_\mu (\not{p} + \not{p}_2) \gamma_{\mu_1} \not{p} \gamma_{\mu_2} (\not{p} + \not{p}_2) \gamma_\nu) \frac{1}{(p + p_2)^4} \right.$$

$$+ \text{tr}(\not{p}_1 \gamma_{\mu_1} (\not{p}_1 + \not{p}_2) \gamma_\mu \not{p} \gamma_\nu (\not{p}_1 + \not{p}_2) \gamma_{\mu_2}) \frac{1}{(p_1 + p_2)^4}$$

$$- \text{tr}(\not{p}_1 \gamma_{\mu_1} (\not{p}_1 + \not{p}_2) \gamma_\mu \not{p} \gamma_{\mu_2} (\not{p} + \not{p}_2) \gamma_\nu) \frac{1}{(p + p_2)^2 (p_1 + p_2)^2}$$

$$\left. - \text{tr}(\not{p}_1 \gamma_\mu (\not{p} + \not{p}_2) \gamma_{\mu_1} \not{p} \gamma_\nu (\not{p}_1 + \not{p}_2) \gamma_{\mu_2}) \frac{1}{(p + p_2)^2 (p_1 + p_2)^2} \right\}.$$

Here, e is the electro-magnetic charge, $C_F = (N_c^2 - 1)/(2N_c)$ arises from the quark color charge and g_s is the strong coupling constant.

The differential phase space for the reaction in Eq.(2) in x is in D -dimensions given by

$$d\text{PS} = \frac{1}{2} \frac{1}{(4\pi)^{D/2-1}} \frac{(Q^2)^{D/2-2}}{\Gamma(D/2-1)} (1-x)^{D/2-2} \int_0^1 dy [y(1-y)]^{D/2-2}, \quad (4)$$

where the relation of the scattering angle θ in the center-of-mass system for q is related to y as $y = \frac{1}{2}(1 + \cos \theta)$.

It is advantageous to automatize the calculation, e.g. with the symbolic computer algebra system FORM, see <https://github.com/vermaseren/form>. A FORM script `nloEpEM.frm` is available on the web-page of this lecture and the output of running the FORM script is listed below.

```

*****
*
* computation of the time-like PTqq0 splitting function in QCD
* run with form (available from https://github.com/vermaseren/form)
*
* S. Moch, 2025
*
* Output definitions:
* den(x) = 1/x;   pow(x,a) = x^a;   InvGamma(x?) = 1/Gamma(x);
*
*****
#-

```

```
nlorm2 =
```

```

+ delta_([1-x])*Gamma(1 - ep)*InvGamma(1 - 2*ep)*cf*ep^(-2) * (
+ 4
)

+ delta_([1-x])*Gamma(1 - ep)*InvGamma(1 - 2*ep)*cf*ep^(-1) * (
+ 4*den(1 - 2*ep)
- 2*den(1 - 2*ep)*den(2 - 2*ep)
)

+ Gamma(1 - ep)*InvGamma(1 - 2*ep)*cf*ep^(-1) * (
- 4*pow([1-x], - 1 - ep)*pow(x, - 2*ep)
+ 2*pow([1-x], - ep)*pow(x, 1 - 2*ep)
+ 2*pow([1-x], - ep)*pow(x, - 2*ep)
);

```

```
0.05 sec out of 0.05 sec
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