
Exploring the Impact of Extra Dimensions on the Equation of State of Baryonic Matter and the Structure of Neutron Stars

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Exploring the Impact of Extra Dimensional spacetime on Neutron Star Structure and Equation of State

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Extra Spatial Dimensions?

- ❑ There are many theoretical considerations that support the existence of extra spatial dimensions ([Lugones & Arbañil 2017](#), [Chakravarti et al. 2018](#), [Arbañil et al. 2019](#))
- ❑ Initially, these extra dimensions naturally emerged in string theory, which mandates ten or more dimensions.
- ❑ Later, the presence of extra dimensions was revisited to address the ‘gauge hierarchy problem’ ([Arkani et al 1998](#), [Antoniadis1988](#), [Perez-Lorezana2005](#)). This problem arises due to the significant and seemingly unrelated hierarchy between the electro-weak symmetry breaking scale (around 10^3GeV) and the Planck scale (about 10^{18}GeV).

Extra Spatial Dimensions?

- ❑ There are a few observational arguments that favor the presence of extra spatial dimensions (GW170817) [Pardo et al. 2018, Abbott et al 2019]

- ❑ GW luminosity distance of 40^{+8}_{-14} Mpc (Abbott et al. 2017)

Vs

EM luminosity distance of $40.7^{+2.4}_{-2.4}$ Mpc (Cantiello et al. 2018)

➡ *Possibility of unequal luminosity distances due to extra spatial dimension !!*

cannot be dismissed

(Deffayet and Menou 2007,
Liu et al. 2023)

ries of gravity with a characteristic length scale R_c of the order of the Hubble radius $R_H \sim 4$ Gpc, such as the well known Dvali-Gabadadze-Porrati (DGP) models of dark energy [88, 89], small transition steepnesses ($n \sim \mathcal{O}(0.1)$) are excluded by the data. Our analysis cannot conclusively rule out DGP models that provide a sufficiently steep transition ($n > 1$) between GR and the onset of gravitational leakage. Future LIGO-Virgo observations of binary neutron star mergers, especially at higher redshifts, have the potential to place stronger constraints on higher-dimensional gravity.

B. P. Abbott et al. Phys. Rev. Lett. 123, 011102 (2019)

M-R curves for Neutron Stars in 4D GR

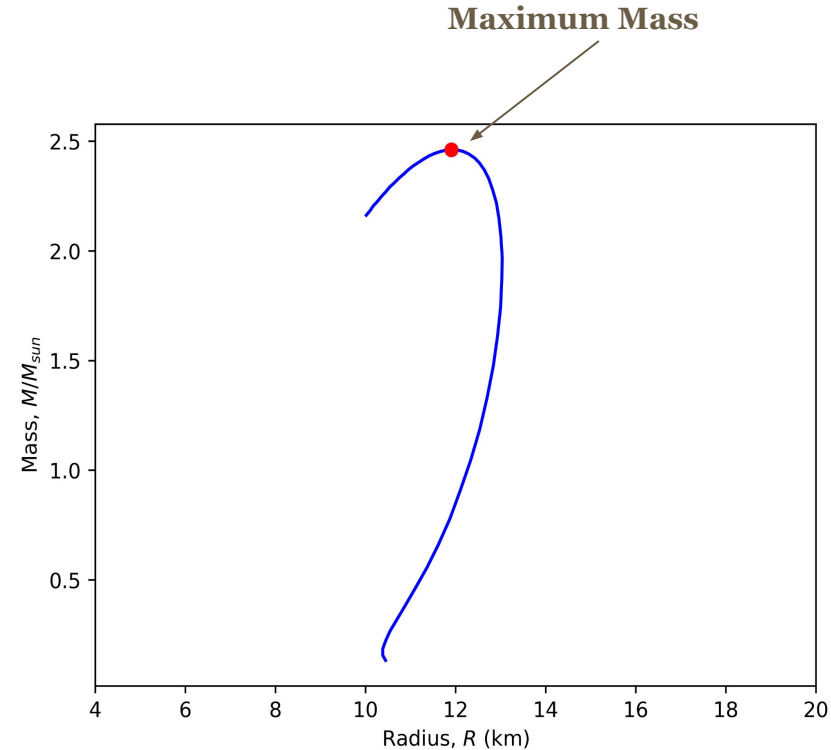
- ❑ The Tolman-Oppenheimer-Volkoff (TOV) equations for a spherically symmetric and static neutron star (NS)

$$\frac{dm}{dr} = 4\pi\rho r^2 \quad - (A) \quad \leq \text{Mass balance Equation}$$

$$\frac{dp}{dr} = -\frac{(\rho + p)(m + 4\pi r^3 p)}{r(r - 2m)} \quad \leq \text{Force Balance Equation}$$

- (B)

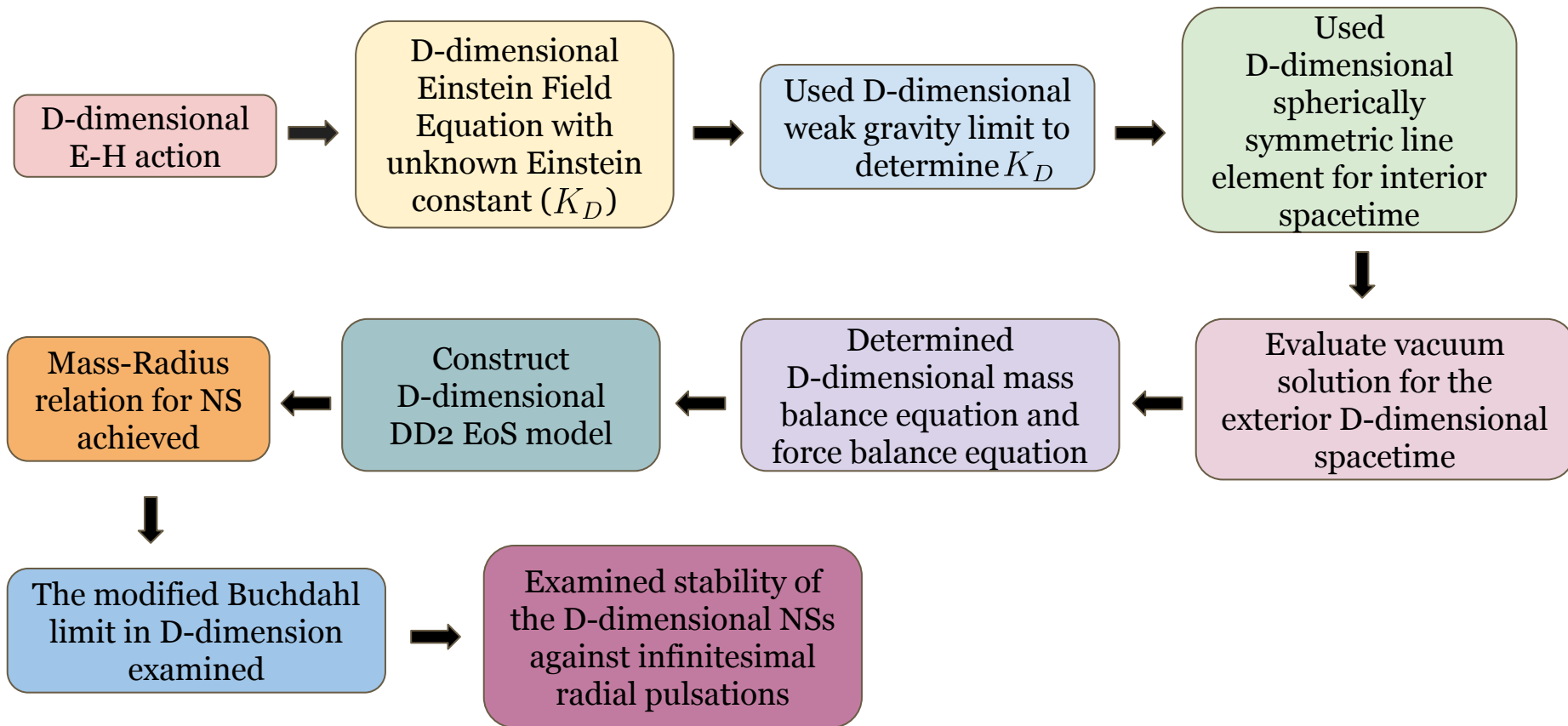
- ❑ Equation of State: To explain the **strong nuclear force** within NSs we have considered realistic **DD2 EoS model**, which describes the nuclear matter, consisting of baryonic components like n, p, alongside leptons such as e and μ . Here we consider baryons interact through the exchange of σ , ω and ρ mesons.



What did we do in <https://arxiv.org/abs/2403.07174>

- ❑ Investigated the influence of higher-dimensional spacetime on both the macroscopic (e.g., mass, radius) and microscopic (e.g., density, pressure) properties of neutron stars (NSs) within a higher-dimensional framework.
- ❑ Extended a realistic nuclear equation of state (EoS) to higher dimensions to accurately describe the strong interactions within neutron matter.
- ❑ Finally, verified the stability of NSs in higher-dimensional spacetime using the Buchdahl condition and stability analysis against infinitesimal radial pulsations.

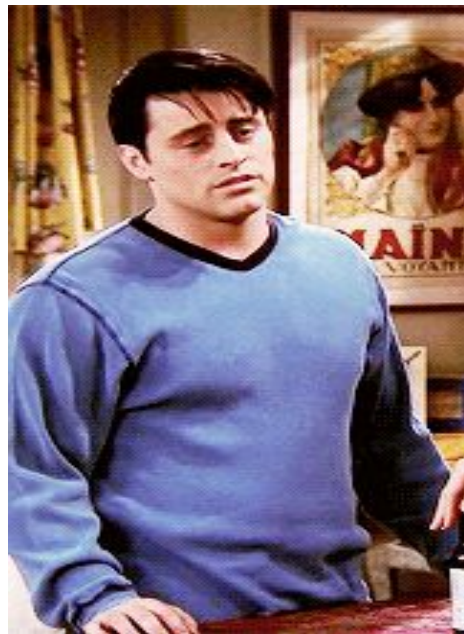
Plan of present work



Challenges

What is the value of Gravitational constant in D-dimension?

What is the value of planck length in D-dimension?



$$\xi_D = (l_c)^{D-4} = G_D/G_4 \Rightarrow l_c \text{ is the length of extra compact dimension}$$

(B Zwiebach (2009), A First Course in String Theory)

D-dimensional Einstein Field Equation

- Generalized Einstein-Hilbert (EH) Action in D-dimensions:

$$S = \frac{1}{K_D} \int d^D x \mathcal{R} \sqrt{-g} + \int d^D x \mathcal{L}_m \sqrt{-g} \quad - (1)$$

$$\text{Einstein's Constant: } K_D = 2 \frac{D-2}{D-3} \frac{G_D}{c^4} S_{D-2}$$

$$\text{Area of Unit Sphere: } S_{D-2} = \frac{2 \pi^{\frac{D-1}{2}}}{\Gamma(\frac{D-1}{2})} \quad - (2)$$

- Stress-Energy Tensor for Isotropic Fluid: $T_{\mu\nu} = \left(\rho_D + \frac{p_D}{c^2} \right) u_\mu u_\nu - p_D g_{\mu\nu} \quad - (4)$

- Generalized Einstein Field Equation: $\mathcal{R}_{\mu\nu} - \frac{1}{2} \mathcal{R} g_{\mu\nu} = \frac{D-2}{D-3} S_{D-2} \frac{G_D}{c^4} T_{\mu\nu} \quad - (5)$

Basic Stellar Equations in D-dimension

- ❑ The exterior Schwarzschild spacetime metric in D-dimensional spacetime

$$ds^2 = \left(1 - \frac{2MG_D}{c^2(D-3)r^{D-3}}\right) dt^2 - \left(1 - \frac{2MG_D}{c^2(D-3)r^{D-3}}\right)^{-1} dr^2 - r^2 \sum_{i=1}^{D-2} \left(\prod_{j=1}^{i-1} \sin^2 \theta_j d\theta_i^2 \right) - (6)$$

- ❑ Stellar Structure Equations:

- Mass Balance Equation: $\frac{dm}{dr} = S_{D-2} \rho_D r^{D-2} \quad - (7)$

- Force Balance Equation: $\frac{dp_D}{dr} = -(\rho_D c^2 + p_D) \frac{G_D [S_{D-2} p_D r^{D-1} + c^2 m (D-3)]}{c^2 r [c^2 (D-3) r^{D-3} - 2m G_D]} \quad - (8)$

- ❑ Simplifying Assumptions:

- Mass, density, and pressure in D-dimensions:

$$\tilde{m} = m G_D / (D-3), \tilde{\rho} = \rho_D G_D, \tilde{p} = p_D G_D \quad - (9)$$

- Set $G_4 = 1$ and $c = 1$ for simplicity.

DD2 EoS model in D-Dimension

- ❑ Relativistic Mean-Field (RMF) Approximation in D-Dimensions is employed to solve Meson Field Equations:

$$m_\sigma^2 \sigma = \sum_{B=n,p} g_{\sigma B} \rho_{DB}^s, \quad - (10)$$

$$m_\omega^2 \omega_0 = \sum_{B=n,p} g_{\omega B} \rho_{DB}, \quad - (11)$$

$$m_\rho^2 \rho_{03} = \sum_{B=n,p} g_{\rho B}^2 \tau_{3B} \rho_{DB}. \quad - (12)$$

- ❑ Baryon Number Density (ρ_{DB}) and Scalar Number Density (ρ_{DB}^s):

$$\rho_{DB}^s = \frac{2J_B + 1}{(2\pi)^{D-1}} S_{D-2} \int_0^{k_B} \frac{k^{D-2} m_B^*}{\sqrt{k^2 + m_B^{*2}}} dk, \quad - (13)$$

$$\rho_{DB} = \frac{2J_B + 1}{(2\pi)^{D-1}} S_{D-2} \int_0^{k_B} k^{D-2} dk. \quad - (14)$$

- ❑ Effective Baryon Mass : $m_B^* = m_B - g_{\sigma B} \sigma$ - (15)

DD2 EoS model in D-Dimension

❑ Dirac Equation for Baryon Fields: $\left[(i\gamma^\mu \partial_\mu - \gamma^0 \Sigma_0^\tau) - m_B^* \right] \psi_B = 0$ - (16)

Here:

$$\Sigma_0^\tau = g_{\omega B} \omega_0 + g_{\rho B} \tau_{3B} \rho_{03} + \Sigma_0^R \quad - (17)$$

Where,

$$\Sigma_0^R = \sum_B [-g_{\sigma B}' \sigma \rho_{BD}^s + g_{\omega B}' \omega_0 \rho_{DB} + g_{\rho B}' \tau_{3B} \rho_{03} \rho_{DB}]$$

(Rearrangement term) - (18)

❑ The energy density (ρ_D) for NS matter

$$\begin{aligned} \rho_D = & \frac{1}{2} m_\sigma^2 \sigma^2 + \frac{1}{2} m_\omega^2 \omega_0^2 + \frac{1}{2} m_\rho^2 \rho_{03}^2 + S_{D-2} \sum_{B=n,p} \frac{2J_B + 1}{(2\pi)^{D-1}} \int_0^{k_B} k^{D-2} \sqrt{k^2 + m_B^{*2}} dk \\ & + S_{D-2} \sum_{l=e,\bar{\mu}} \frac{2J_l + 1}{(2\pi)^{D-1}} \int_0^{k_l} k^{D-2} \sqrt{k^2 + m_l^2} dk, \end{aligned} \quad - (19)$$

❑ the pressure (p_D)

$$\begin{aligned} p_D = & -\frac{1}{2} m_\sigma^2 \sigma^2 + \frac{1}{2} m_\omega^2 \omega_0^2 + \frac{1}{2} m_\rho^2 \rho_{03}^2 + \Sigma_0^R \sum_{B=n,p} \rho_{DB} \\ & + \frac{S_{D-2}}{D-1} \sum_{B=n,p} \frac{2J_B + 1}{(2\pi)^{D-1}} \int_0^{k_B} \frac{2(D-2)}{\sqrt{k^2 + m_B^{*2}}} dk + \frac{S_{D-2}}{D-1} \sum_{l=e,\bar{\mu}} \frac{2J_l + 1}{(2\pi)^{D-1}} \int_0^{k_l} \frac{2(D-2)}{\sqrt{k^2 + m_l^2}} dk \end{aligned} \quad - (20)$$

DD2 EoS model in D-Dimension

The nucleon-meson density-dependent couplings (Typel et al 1999, 2005, 2010) :

$$g_{\alpha B}(n) = g_{\alpha B}(n_0) f_{\alpha}(x) \quad \text{where } x = n/n_0$$

n = total baryon density
 n_0 = saturation density

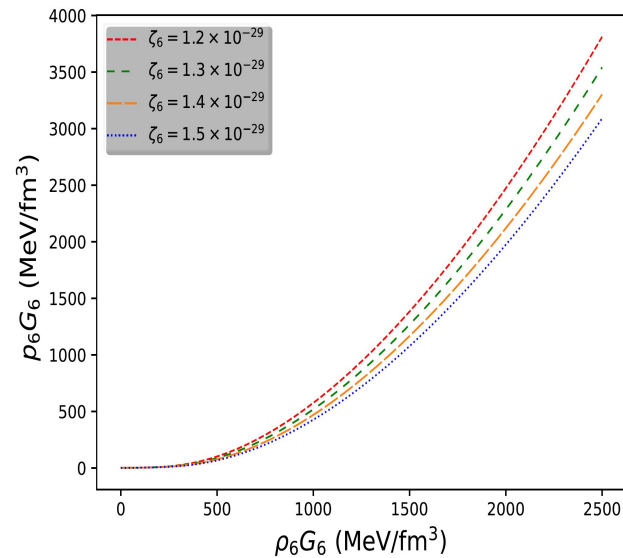
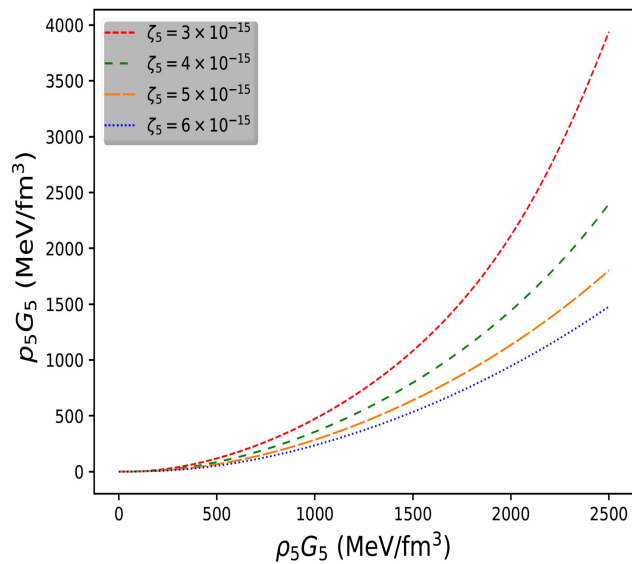
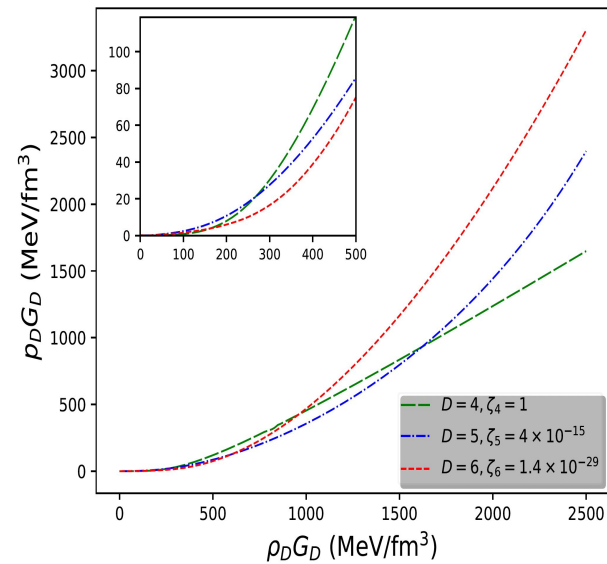
- (21)

For σ and ω meson: $f_{\alpha}(x) = a_{\alpha} \frac{1 + b_{\alpha}(x + d_{\alpha})^2}{1 + c_{\alpha}(x + d_{\alpha})^2}$ - (22)

For ρ meson: $f_{\alpha}(x) = \exp[-a_{\alpha}(x - 1)]$ - (23)

The values of the parameters $g_{\alpha B}(n_0)$, a_{α} , b_{α} , c_{α} , and d_{α} for $\alpha = \sigma, \omega$ and ρ are obtained through fitting the nuclear properties.

DD2 EoS model in D-Dimension



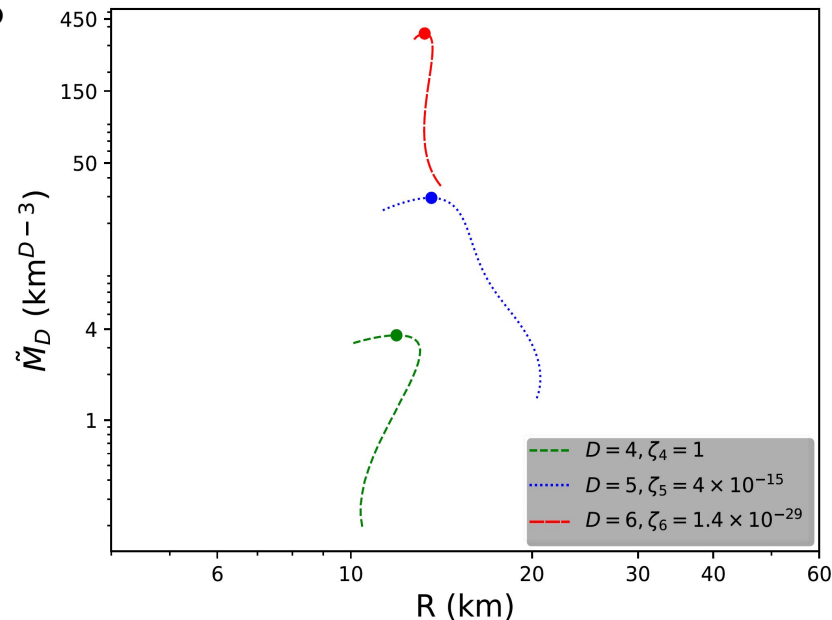
Mass-radius curve for Neutron Stars in D-dimensions

EH action => Einstein field Equation => Schwarzschild exterior metric=> mass balance equation => From EFE to TOV equation => modified EOS => MR curve $\longleftarrow$ Need to explore all in D-dimensional spacetime

Variation of the total mass ($\log(MG_D/(D-3))$) in $\text{km}^{(D-3)}$ with the total radius R in km for the different D-dimensional cases, such as D=4, 5 and 6

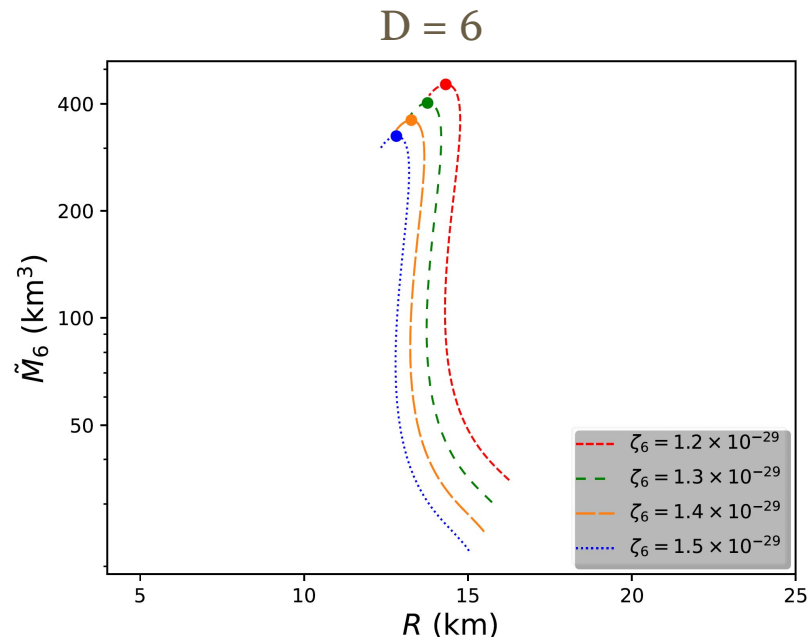
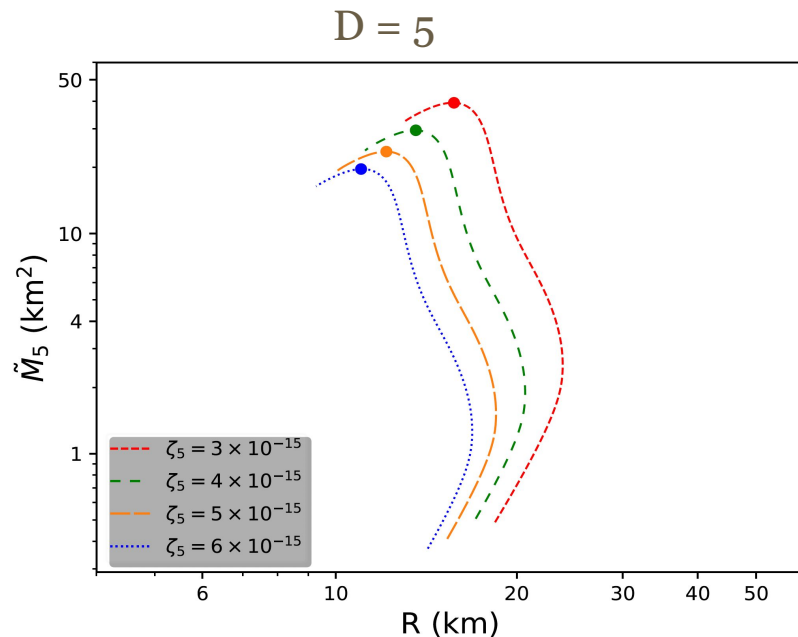
- ❑ $G_4 \neq G_5 \neq G_6 \dots \neq G_D$,
we don't know the values for G_D
- ❑ We can employ dimensional analysis to introduce a scaling parameter ζ_D given by

$$G_D/G_4 = \zeta_D \text{ km}^{D-4} \quad - \text{ (C)}$$



Mass-radius curve for neutron stars in D-dimension

Variation of the total mass ($\log(MG_D/(D-3))$) in $\text{km}^{(D-3)}$ with the total radius R in km for the $D=5$ and 6 dimensional cases:

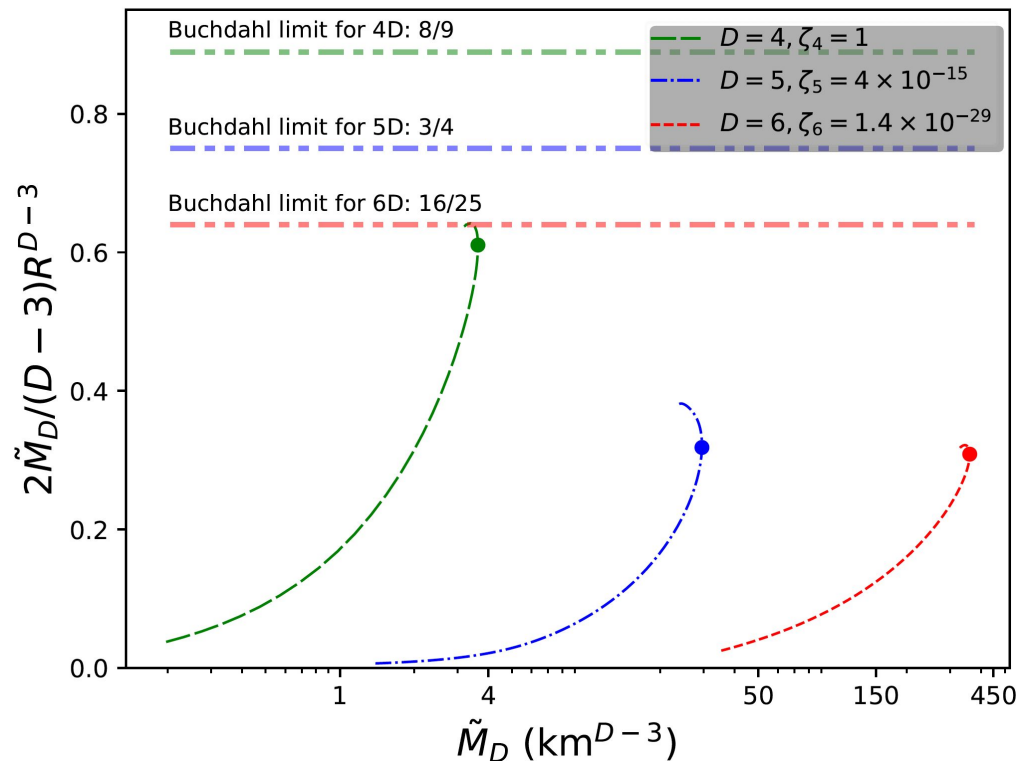


Buchdahl Limit and Surface Redshift for NSs in D-dimension

Buchdahl limit for D-dimension spacetime:

$$\frac{MG_D}{c^2(D-3)R^{D-3}} < \frac{2(D-2)}{(D-1)^2} \quad - (D)$$

Buchdahl limit ensures that a neutron star is **stable against gravitational collapse** and provides an **upper limit on how compact** a neutron star can be without collapsing into a black hole

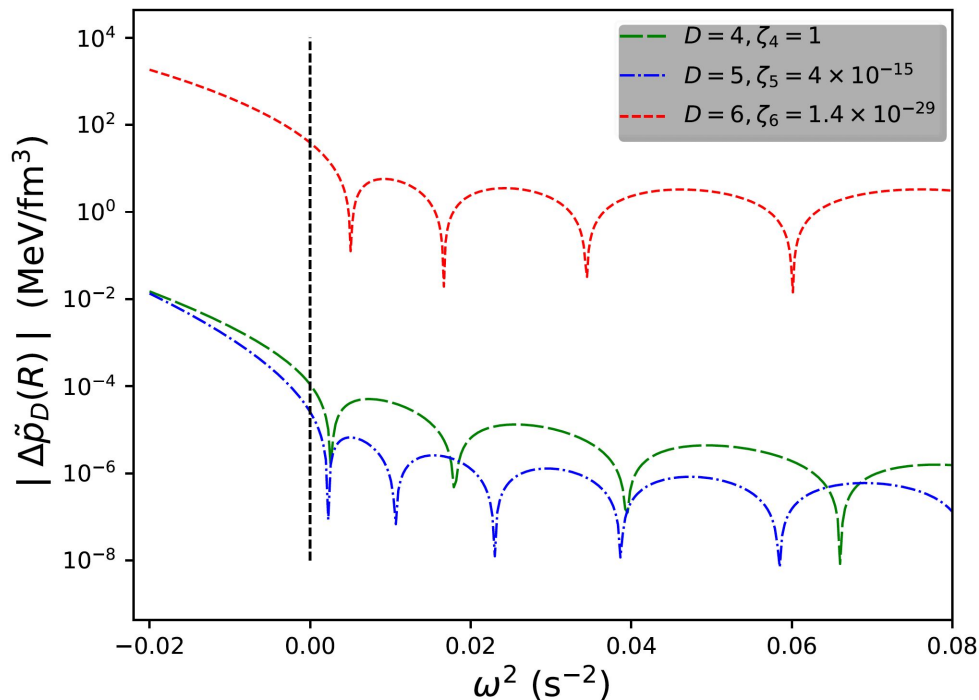
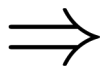


Testing stability against infinitesimal Radial Pulsation

Lagrangian perturbation of the radial pressure ($\Delta\tilde{P}(R) = \Delta P(R)G_D$) for a range of test values ω^2 within different D-dimensional spacetimes

ω = eigenfrequency

Lagrangian perturbation
of the radial pressure



$\omega^2 > 0$ ensures neutron star is stable against radial perturbations for $D = 4, 5$ and 6 .

Numerical values of some physical parameters

Value of D	Value of Maximum Mass $MG_D/(D-3)$ in $\text{km}^{(D-3)}$	Value of Corresponding Radius (km)	Central Density $\tilde{\rho}_{D,c}$ (MeV fm^{-3})	Central Pressure $\tilde{p}_c$ (MeV fm^{-3})	$\frac{2MG_D}{(D-3)R^{D-3}}$	Surface Redshift	Compactified Maximum Mass in 4-dimension ($M_\odot$)
4	3.63	11.91	1064.85	501.78	0.61	0.60	2.46
5	29.49	13.58	1459.85	754.40	0.32	0.21	2.50
6	360.25	13.27	2291.54	2782.94	0.31	0.20	3.31

Value of ζ_5	Value of Maximum Mass $MG_D/(D-3)$ in $\text{km}^{(D-3)}$	Value of Corresponding Radius (km)	Central Density $\tilde{\rho}_{D,c}$ (MeV fm^{-3})	Central Pressure $\tilde{p}_c$ (MeV fm^{-3})	$\frac{2MG_D}{(D-3)R^{D-3}}$	Surface Redshift	Compactified Maximum Mass in 4-dimension ($M_\odot$)
3×10^{-15}	39.31	15.72	1077.23	547.73	0.32	0.21	2.88
4×10^{-15}	29.49	13.58	1459.85	754.40	0.32	0.21	2.50
5×10^{-15}	23.59	12.13	1841.52	960.35	0.32	0.21	2.24
6×10^{-15}	19.66	11.01	2296.87	1245.24	0.32	0.22	2.05

A brief Summary

- ❑ Higher dimensions significantly affect the structural properties of NSs, including density and pressure profiles, as well as enclosed masses, due to the modified DD2 model EoS.
- ❑ The mass-radius analysis shows a progressive increase in total mass for a constant central density as the dimensionality increases, maintaining the stability criterion
$$dM/d\rho_c > 0$$
- ❑ Higher-dimensional NSs exhibit stability against radial oscillations, adhering to the modified Buchdahl limit for higher dimensions.
- ❑ *We are exploring the potential applications how to match with any existing observational signatures*



Thank you !

Video credit: Interstellar movie