The LHC: always exciting

Kajari Mazumdar

Department of High Energy Physics Tata Institute of Fundamental Research Mumbai.

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Plan of the talk

- Introduction
- Physics platter:

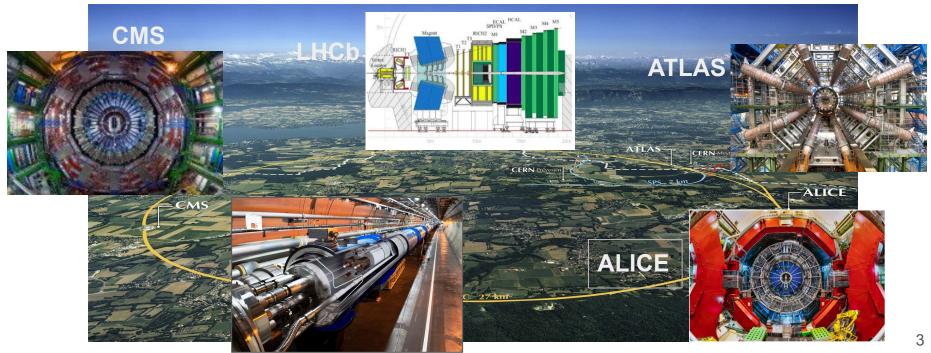
Measurement of the W-mass Search for multi-Higgs productions Top quark front A brief from the heavy ion sector Another from the heavy flavor domain

- Beyond the LHC?
- Conclusion

The large hadron collider (LHC)

- Conceived in early 1980s
- Start of the LHC construction: 1998
- First p-p collision: November, 2009
- Higgs boson discovery: 2012 first motivation for the project
- To operate for ~2 more decades!

Motivation 2: *unfinished job as yet.* No direct evidence for physics beyond the Standard Model (SM)



Proton-on-proton collisions at the LHC

Excellent performance of the LHC machine for more than a decade of operation.

$$N_{\text{observed}} = \sigma_{\text{process}} \times \epsilon \times A \times (Ldt) + N_{\text{background}}$$

p-p collision data collected by each of the ATLAS and CMS experiments:

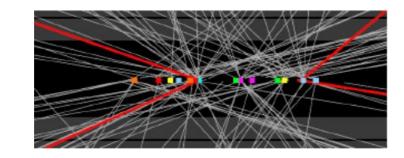
- Run-1 at √s = 7, 8 TeV (2010 2012) : *L* ~ 20 fb⁻¹
- Run-2 at √s = 13 TeV (2015 2018): *L* ~ 140 fb⁻¹
- Run-3 at \sqrt{s} = 13.6 TeV (2022+2023+2024): \mathcal{L} ~ 105 fb⁻¹

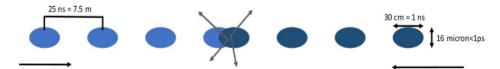
In 2024, peak instantaneous luminosity L ~ 2.5X10³⁴ cm⁻²s⁻¹

- ~ 65 p-p collisions/crossing on average
- \rightarrow pile up for a triggered event

Run-3 will continue till 2026, expected total \mathcal{L} ~ 250 fb ⁻¹

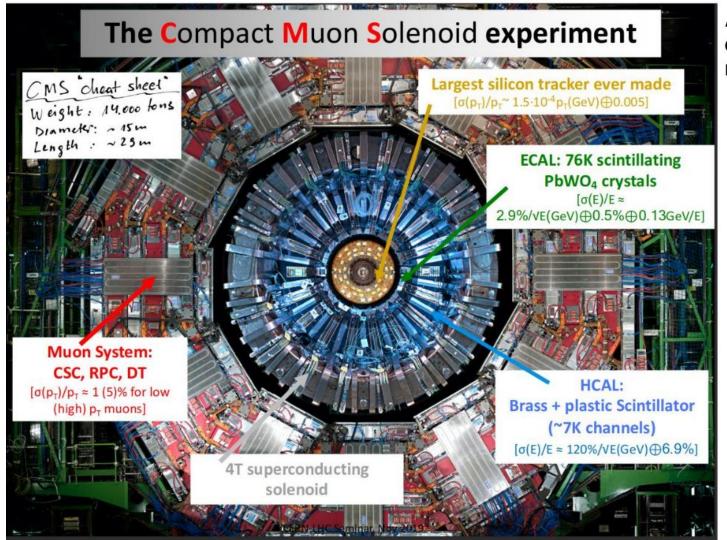
Low luminosity datasets in 2017 $\mathcal{L} = 298 \text{ pb}^{-1} @ \sqrt{s} = 5.02 \text{ TeV}$ $\mathcal{L} = 201 \text{ pb}^{-1} @ \sqrt{s} = 13 \text{ TeV}$ + Heavy ion runs



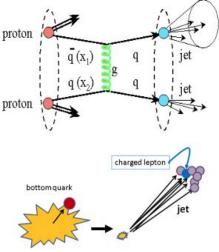


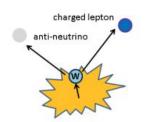
Pileup: multiple interactions at **different** hard-scales

Data collected by each experiment ~ 150 PB



ATLAS and CMS experiments are THE multi-messenger tools!



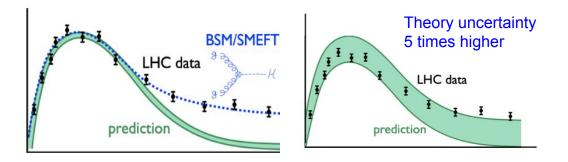


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Dividend of collective efforts: achieve more than anticipated!

- Huge amount of physics harvest including Higgs boson discovery, measurements, search for rare processes and physics beyond the SM.
- Enormous progress in computing capabilities.
- Precision theoretical description of the crucial processes
- Extensive understanding of the performance of the experiments.
- Judicious application of sophisticated machine learning techniques!
 - eg. b-jet identification: Deep Jet in CMS: <u>JINST 15 (2020) P12012</u>, <u>CMS-DP-2021-014</u>

We cannot afford to miss a discovery!



Note: LHC is the only APPROVED future project of HEP as of today

 \Rightarrow maximise the physics harvest.

Electroweak physics

Milestones

- Neutral currents: 51 yrs
- QCD: 51 yrs

....

- W, Z turns 41
- Top: 29
- BEH particle: 12
- ElectroWeak sector one of the most successful and well tested parts of the Standard Model (SM).
- Precision measurements have guided discovery of the top quark, Higgs boson using accurate prediction for electroweak observables.

 M_w^2

- Input parameters for predictions are all known now!
 - Z Mass
 - Top Mass
 - Higgs Mass

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⇒ over-constrain the SM.

$$\left(1 - \frac{M_W^2}{M_Z^2}\right) = \frac{\pi \,\alpha_{\rm EM}(M_Z)}{\sqrt{2}G_F(1 - \Delta r)}$$
Higher order corrections, depend on m_t, m_H, m_{BSM}?

Theoretical uncertainty of $\Delta M_{top} = 2.1 \text{ GeV} \Rightarrow \Delta M_{W} = 1.9 \text{ GeV}$ Today LHC: $\Delta m_{t} = 330 \text{ MeV}$ $\Delta m_{H} = 110 \text{ MeV}$

Nature 633, 745-746 (2024)

Precision measurements at the LHC crucial to get hint about heavier particles not directly accessible now.

Measurement of W mass

 $m_W^{\rm SM} = 80355 \pm 6 \,\,{
m MeV}$

• Humongous effort by theory and experimental communities to push the measured ΔM_w to < 10 MeV

LEP measurement: ΔM_7 = 2 MeV, precision of 22 parts in a million LHCb

- Complex, indirect m_w measurement
 - \rightarrow requires $\mathcal{O}(5-7)$ years

 \rightarrow Only four W-boson mass measurements in the last 10 years.

 $2021 - LHCb m_{W} = 80.354 \pm 0.032 \text{ GeV}$

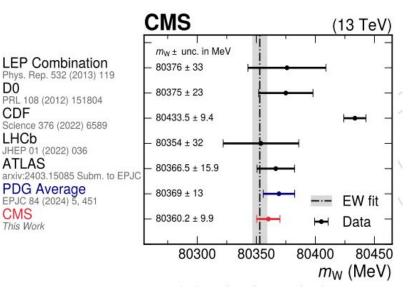
2022 – <u>CDF</u> m_w = 80.4335 ± 0.0094 GeV, 4.2 M events

2024 – <u>ATLAS</u> m_w = 80.3665 ± 0.0159 GeV [9.8 (stat) * 12.5 (syst) MeV] (14 M events, using 7 TeV data 4.6 /fb, PDF CT18)

2024 – <u>CMS</u> m_w = 80.3602 ± 0.0099 GeV [2.4 (stat) * 9.6 (syst) MeV] (1 B events, using ~ 10% of 13 TeV data, uses CT18Z PDF)

Used 5 billion Monte Carlo samples

History of m_w measurements



Uncertainty due to muon momentum scale ATLAS: ΔM_W ~ 6 MeV CMS: ΔM_W ~ 4.4 MeV CDF: ΔM_W ~ 3 MeV

Uncertainty due to parton density function ATLAS: $\Delta M_W \sim 5 MeV$ CMS: $\Delta M_W \sim 2.8 MeV$ CDF: $\Delta M_W \sim 3.9 MeV$

W boson production and decay

- p-pbar collisions: W bosons mostly produced in the same helicity state.
- p-p collisions: both positive and negative polarization states.
- Large PDF-induced W-polarisation uncertainty affecting the ${\rm p_T}$ lepton distribution, peak at ${\rm m_W}/2$

 $\vec{p}_{\mathrm{T}}^{\mathrm{miss}} = -\left(\vec{p}_{\mathrm{T}}^{\ell} + \vec{u}_{\mathrm{T}}\right)$

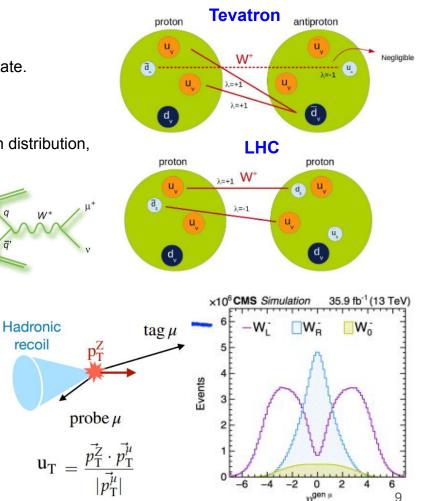
• Sea-quark PDFs play a larger role at the LHC.

$$m_{\rm T}^{\rm W} = \sqrt{2p_{\rm T}^{\ell}p_{\rm T}^{\rm miss}(1 - \cos\Delta\phi_{\ell\vec{p}_{\rm T}^{\rm miss}})}$$

Additional QCD complications

- 1. Heavy-flavour-initiated processes
- 2. W+, W- and Z are produced by different light flavour fractions
- 3. Larger gluon-induced W production

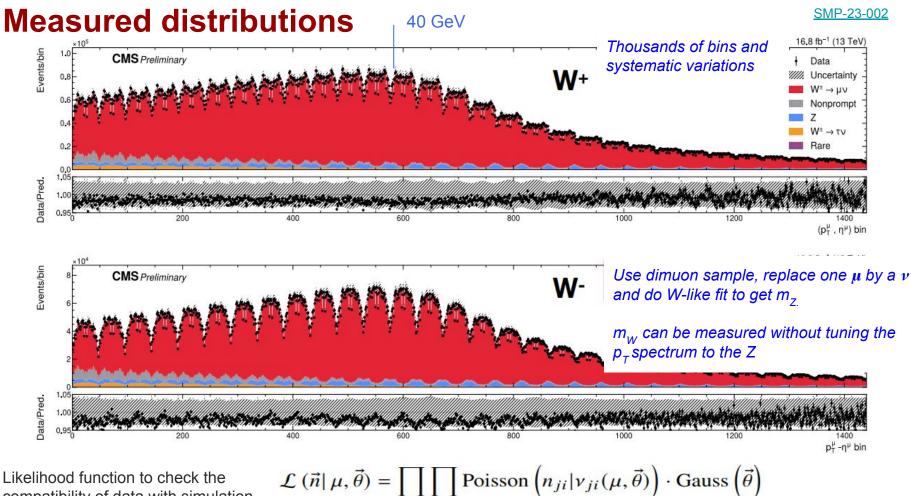
Note: CDF and LHC experiments performed at different times, with different baseline PDFs and QCD tools, different experimental conditions,..



CMS strategy: highlights

- Measure m_W by fitting the observed 3D distribution of lepton charge, η^{ℓ} and $p_T^{-\ell}$ using template shapes for the signal and the background processes.
- Used large data volume collected during Run2.
- High pileup (avg.~ 45) data makes p_{τ}^{miss} resolution worse extracting m_{W} from m_{τ}^{W} particularly challenging
- Focus on lepton kinematics, restrict to muon channel (better calibration).
- Data has strong constraining power \Rightarrow utlise sensitivity to PDF from η^{ℓ} distribution
- Low p_T^W region treated in the best possible way (N3LL) + analytical resummation in recoil, high $p_T^W @ N2LO$
- Critical to control the lepton efficiency and momentum calibration. ($\Delta p_T^{\mu} \sim 10^{-4} \Rightarrow \Delta m_W \sim 8 \text{ MeV}$)
- p_T^μ distribution strongly dependent on the theoretical modeling of p_T^W distribution

 → calibration of the p_T^μ measurement only uses muons from the J/ψ → μμ resonance.
 → sample of muons from Y(1S) and Z decays used for an independent validation of the p_T^μ calibration
- 100 M selected (out of 300 M data) W→µv events 87% signal rest mainly QCD multijet events: contributions estimated in data-driven method
- 7 M selected $Z \rightarrow \mu \mu$ events m_{$\ell\ell$} in [60, 120] GeV >99% signal Validation:measured m₇ = 91 182 ± 7 (stat) ± 12 (syst) = 91 182 ± 14 MeV ,

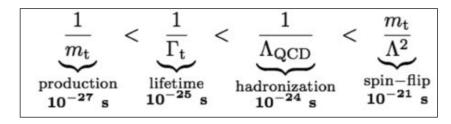


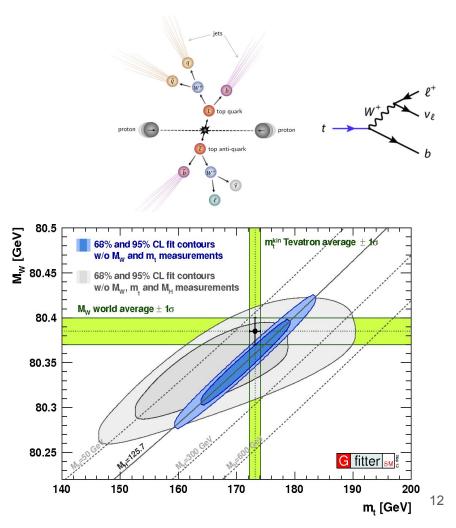
compatibility of data with simulation

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Physics of the top quark

- Heaviest known elementary particle so far.
- Extremely short-lived, decays before hadronization.
 ⇒ spin information transferred to decay products.
 - \rightarrow possible to observe properties of the bare quark
- Abundant production at the LHC (in Run 2 ~10⁸ per expt.)
- m_t = 172.52 ± 0.33 [0.14 (stat) ⊕ 0.33 (syst)] GeV
 → impressive precision of 0.18% 2403.01313
- Mass/ Yukawa coupling: free parameter in SM, \rightarrow to be measured in experiments: y_{t} = $\sqrt{2}~m_{t}/\textit{v}$.
- **m**_t key factor in the stability of SM Higgs potential and hence the evolution of the universe.





Spin entanglement of top quark pairs

- Entanglement: key feature of QM
- Spin correlation explored typically in low energy regime in various context.
- LHC: a pair of quarks available at relativistic energies: minimum 2-qubit system.
- Challenge for observation of spin-entanglement in tt system.
 - \Rightarrow Cannot control the internal degrees of freedom in the initial state.
 - \Rightarrow Entanglement can be detected only with dedicated analysis in a restricted phase space.
- Study two-qubit states of ttbar at production threshold ⇒ mostly singlet → rotational invariance with well-specified fiducial phase-space → maximally correlated or entangled!

Spin density matrix of 2-qubit system

$$\rho = \frac{1}{4} \left[I_4 + \sum_i (B_i^+ \sigma^i \otimes I_2 + B_i^- I_2 \otimes \sigma^i) + \sum_{i,j} C_{ij} \sigma^i \otimes \sigma^j \right].$$

 I_n : normalization B_i : individual polarization C_{ij} : spin correlation

each diagonal element of C corresponds to the spin correlation in a particular direction, \rightarrow good entanglement witness.

Require Tr(C) + 1 < 0

Nature 633 (2024) 542

Observation of spin entanglement at the highest energy

1	$\mathrm{d}\sigma$	$1 + \mathbf{B}^+ \cdot \hat{\mathbf{q}}_+ - \mathbf{B}^- \cdot \hat{\mathbf{q}} \hat{\mathbf{q}}_+ \cdot \mathbf{C} \cdot \hat{\mathbf{q}}$
$\overline{\sigma} d$	$\Omega_+ d\Omega$	$=$ (4 π) ²

Entanglement marker $D = -3\cos \varphi$

B : top polarization q: individual momentum

 $D = \frac{\operatorname{Tr}\left[\mathbb{C}\right]}{3} \Rightarrow D < -\frac{1}{3}$

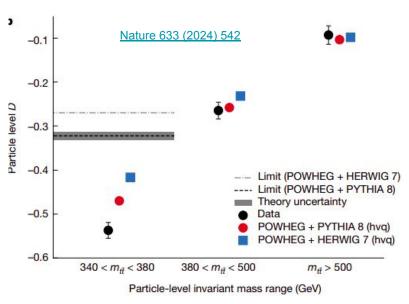
 \rightarrow observable dependent on the angle ϕ between the charged leptons in the rest frame of their parents.

 \rightarrow cos ϕ can be measured experimentally in an ensemble dataset.

ATLAS: D= -0.547 \pm 0.002 (stat) \pm 0.020 (syst) for 340 < m_t < 380 GeV

Significance more than 5σ compared to null hypothesis of no-entanglement

CMS also confirms observation of spin-entanglement Analysis takes into account parton level entanglement combined with non-relativistic effects, contribution from toponium production.



2406.03976 Submitted to RPP

Observation of 4 top production

Heaviest final state : σ(tttt) ~ 13.4 ±1.4 fb
 @NLO QCD, NLO EWK + NNLL

O(100M) tt events and O(1k) tt tt events, 100 ttt events

Observed (expected) significance 6.1 (4.7) s.d. : ATLAS 5.6 (4.9) s.d.: CMS

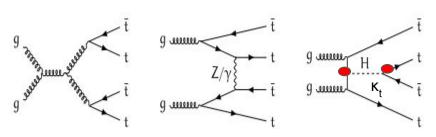
EPJC 83 (2023) 496 PLB 847 (2023) 138290

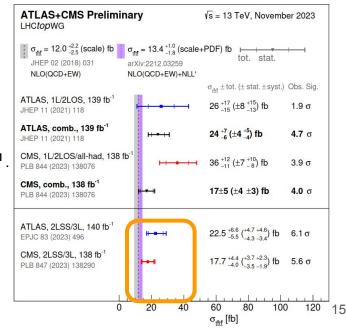
- Possible enhancement in BSM: \Rightarrow Several constraints on new physics.
- Cross section sensitive to top Yukawa coupling, CP properties of $\kappa_t = y_t / y^{SM}$.

Assuming a pure CP-even coupling ($\alpha = 0$), observed upper limit on $|\kappa_{\star}| = 1.9$ at 95% CL

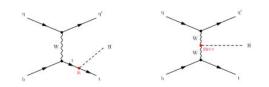
• EFT: constrain 4-fermion interactions of dim-6.

EPJC 84 (2024) 156





Watch out on tH, tWZ processes



Current upper limit on tH production rate ~ 15*SM

tHq, tHW: interference between Higgs boson emission from t or W: in SM, almost maximally destructive \Rightarrow very low rates.

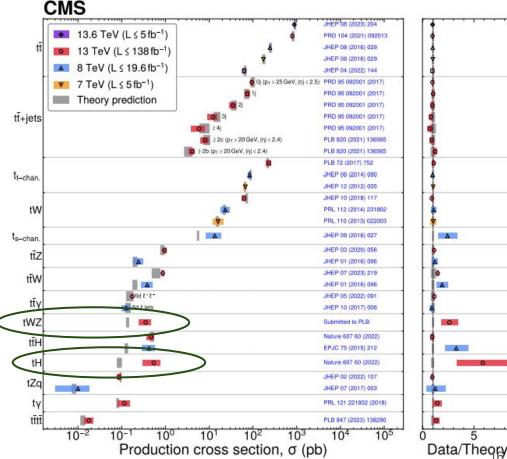
 \Rightarrow BSM can enhance it!

• Sign flip can make $\sigma(tH) >> \sigma(ttH)$

 CP nature of coupling: the relative sign of ttH coupling wrt HWW coupling

Measured σ (tWZ) = 354 ± 54 (stat) ± 95 (syst) fb statistical significance 3.4

 2σ above the SM prediction of 136 ± 9 fb at NLO(QCD)



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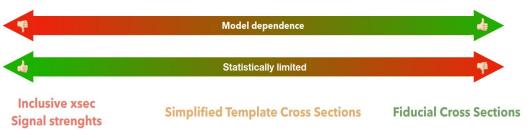
Higgs boson: more than a decade after the discovery

- Total no. of Higgs bosons already produced at each interaction point ~ 15 M
- Higgs signal strength measured with ~ 6% precision, uncertainty still dominated by statistics.
- Higgs physics is now the tool for probing anomalies in data wrt Standard Model predictions or models beyond SM.

Essentially, we study the dynamics of production and decay to understand the properties

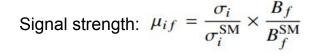
- determine the quantum numbers: charge, spin, parity, mass and width
- check if all the predicted decay channels exist or not
- · look for indications for anomalous couplings
- estimate the shape of the Higgs potential: constrain the trilinear and quartic couplings
- study if the Higgs sector is minimal or an extended one

Evolution in interpretation of data for Higgs physics



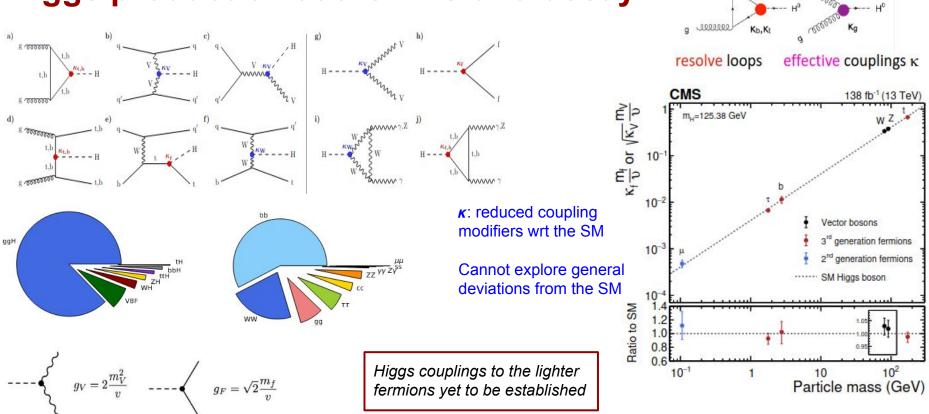


new interaction of Nature



$$(\sigma_i \times B_f) = k_i^2 \sigma_i^{SM} \frac{k_f^2 \Gamma_f^{SM}}{k_H^2 \Gamma_H^{SM}}$$

Higgs production at the LHC and decay



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Properties of the discovered BEH particle compatible with those envisaged for a spin-0 boson in the SM 18

Measurement of Higgs self-interaction

After the electroweak symmetry breaking, Higgs potential:

$$V(h) \simeq \frac{1}{2}m_H^2 h^2 + \lambda h^3 + \frac{1}{4}\lambda h^4 + \dots$$

Determination of Higgs self-coupling parameter λ :

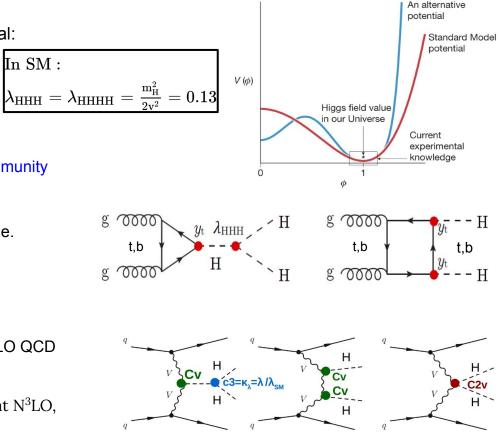
- \rightarrow Currently, THE most important mandate of the community
- \Rightarrow shape of the Higgs potential near the minimum

 \rightarrow related to the evolution of the Universe at the EW scale.

• Inclusive Higgs pair production at the LHC \rightarrow direct access to HHH and VVHH vertices $\rightarrow \kappa_{a}, \kappa_{2V}$

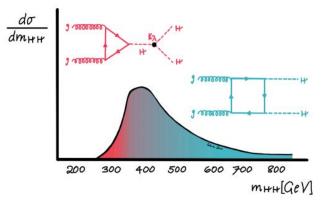
Gluon-gluon fusion $\sigma(gg \rightarrow HH+X) \sim 31$ fb @ 13 TeV with N2LO QCD almost 10³ times smaller than $\sigma(pp \rightarrow H+X)$

Vector Boson fusion: the sub-lead mode, $\sigma_{VBFHH} = 1.73$ fb at N³LO, Unique access to HHVV vertex.

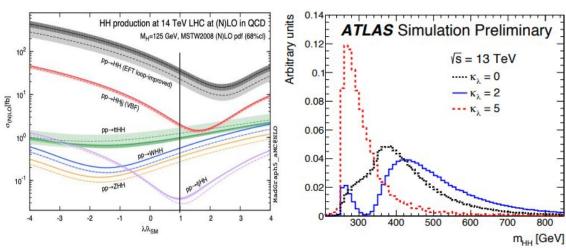


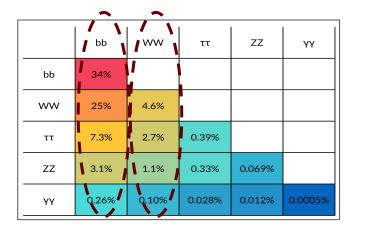
Higgs pair production process is yet to be observed experimentally.

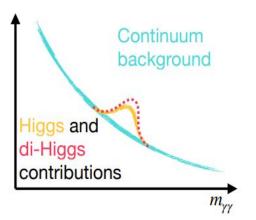
General strategy to search for Higgs boson pair production



Interference among relevant diagrams \rightarrow cross section dependency on the coupling modifiers

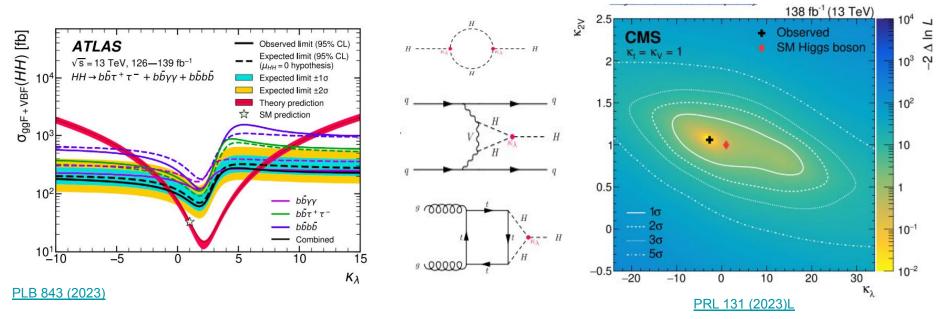






- There is no clear golden channel /combination several promising combination for experimental signatures:
- H→ bb : large branching fraction, large background
- $H \rightarrow \gamma \gamma$ good mass resolution
- $H \rightarrow \tau \tau$:: lower background

Results from search for non-resonant HH production

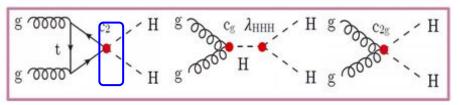


Observed (expected) value of HH production cross section lower than 3.4 (2.5) * SM value Constraints on κ_{λ} with 95% CL : [-1.24, 6.49]

Hypothesis of κ_{2v} (C_{2v}) = 0 excluded with a significance of 6.3 standard deviations.

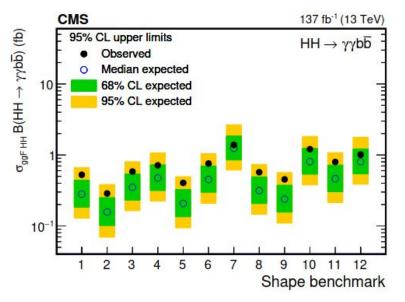
We may have good news on diHiggs measurement sooner than expected!

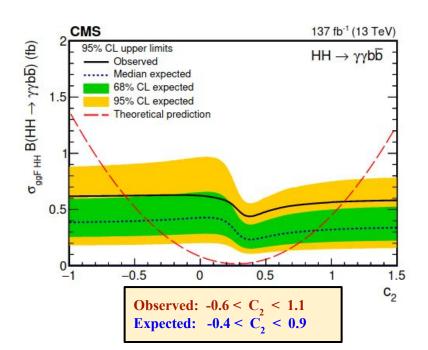
Search for anomalous self-couplings



EFT approach for $gg \rightarrow$ HH process includes **three** types of contact interactions described by dim-6 operators. \rightarrow additional couplings for $gg \rightarrow$ HH compared to SM.

Compatibility of data with different representative BSM scenarios





Resonant HH production

HH final state provides 2-way probe

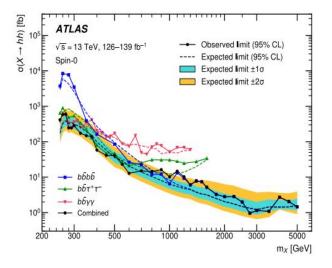
- nature of couplings to BSM particles
- new topologies

Spin-0: Two-Higgs-Doublet-Models completed by an Electroweak Singlet.



The different searches often complementary for different mass ranges.

Results presented in a model agnostic way and often reinterpreted in the 2HDM and MSSM model.



ATLAS: small excess with combined local (global) significance of 3.2 (2.1) at 1.1 TeV.

CMS: limits below 320 GeV and above 1 TeV.

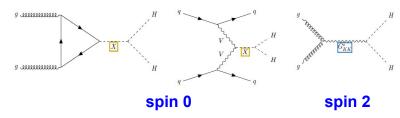
Watch out!!

2405.20040

No significant excess in resonant VBF search.

ATLAS analysis excludes parameter space in the region $2 \le \tan \beta \le 5$, which is not excluded by the standard SUSY H searches

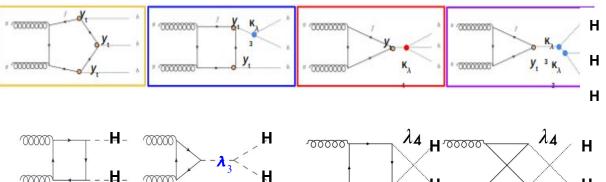
ATLAS-CONF-2024-006



More the merrier!

Can LHC make triple Higgs final state? YES! $\sigma_{SM} = 0.0893$ fb

 $V(h) = \frac{1}{2}m_{H}^{2}h^{2} + \lambda_{3}\nu h^{3} + \frac{1}{4}\lambda_{4}h^{4}$ In BSM, λ_2 can be different from λ_4



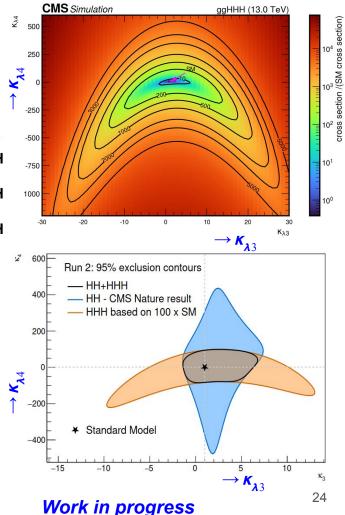
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Multi-Higgs production modes directly accessible at the LHC, mainly via

Di-Higgs production sensitive to λ_{3}

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- Triple Higgs production sensitive to both λ_3 and λ_4 , with stronger dependence on λ_{2}
- Consistent determination of the Higgs potential require combined measurement of λ_3 and λ_4 .



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Particle production at the LHC and stairways to heaven!

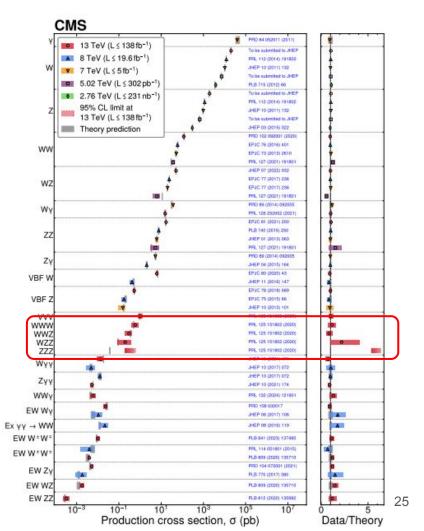
2405.18661 Submitted to Phys.Rep.

Bottom quark: 1 in a hundred pp collisions W boson: 1 in a half a million Z boson: 1 in a million H boson: 1 in a billion HH pair: 1 in a trillion Triple H: 1 in 100 trillion

Measured cross sections varies over 12-13 orders of magnitude

Watch out on triboson productions: WZZ, ZZZ

Measured σ (tWZ) = 354 ± 54 (stat) ± 95 (syst) fb statistical significance 3.4 s.d. (expected: 1.4 s.d.) **2\sigma above the SM prediction** of 136 ± 9 fb at NLO(QCD)



Where are others?

ATLAS SUSY Searches* - 95% CL Lower Limits

ATLAS Preliminary $\sqrt{s} = 13 \text{ TeV}$

		Model	S	ignature	e f£	dt [fb ⁻¹	Mass limit		Reference
	səu	$\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_1^0$	0 <i>e</i> , μ mono-jet 0 <i>e</i> , μ		E_T^{miss} E_T^{miss} E_T^{miss}	140 140 140	7 [1x, 8x Degen.] <u>1.0</u> [8x Degen.] 0.9	1.85 m(χ ⁰)<400 GeV m(∂)-π(X ⁰)=5 GeV 2.3 m(X ⁰)=0 GeV	2010.14293 2102.10874 2010.14293
	Searches	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q \tilde{q} \tilde{\chi}_1^0$			1		Forbidden	1.15-1.95 m($\tilde{\chi}_1^0$)=1000 GeV	2010.14293
		$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\bar{q}W\tilde{\chi}_1^0$	1 e,μ ee,μμ	2-6 jets 2 jets	r miss	140 140	Ś.	2.2 m($\tilde{\chi}_1^0$)<600 GeV 2.2 m($\tilde{\chi}_1^0$)<700 GeV	2101.01629 2204.13072
	sive	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\bar{q}(\ell \ell)\tilde{\chi}_{1}^{0}$ $\tilde{g}\tilde{g}, \tilde{g} \rightarrow qqWZ\tilde{\chi}_{1}^{0}$	0 e,µ	7-11 jets	E_T^{miss} E_T^{miss}	140	ş	1.97 m(\tilde{k}_1^0) <600 GeV	2008.06032
	Inclusive	88, 8-444 MZA 1	SS e, µ	6 jets	D_T	140		.15 m(\tilde{g})-m($\tilde{\chi}_1^0$)=200 GeV	2307.01094
	Inc	$\tilde{g}\tilde{g}, \; \tilde{g} \rightarrow t \tilde{t} \tilde{\chi}_1^0$	0-1 <i>e</i> , μ SS <i>e</i> , μ	3 b 6 jets	$E_T^{\rm miss}$	140 140		2.45 m($\tilde{\chi}_1^0$)<500 GeV 1.25 m($\tilde{\chi}_1^0$)=300 GeV	2211.08028 1909.08457
		$\tilde{b}_1 \tilde{b}_1$	0 <i>e</i> , <i>µ</i>	2 b	$E_T^{\rm miss}$	140	5, 0.68	1.255 m($\tilde{\chi}_{1}^{0}$)<400 GeV 10 GeV<∆m($\tilde{b}_{1}\tilde{\chi}_{1}^{0}$)<20 GeV	2101.12527 2101.12527
	ks on	$\tilde{b}_1\tilde{b}_1,\tilde{b}_1{\rightarrow}b\tilde{\chi}^0_2{\rightarrow}bh\tilde{\chi}^0_1$	0 e,μ 2 τ	6 b 2 b	E_T^{miss} E_T^{miss}	140 140		.23-1.35 $\Delta m(\tilde{\chi}_{2}^{0},\tilde{\chi}_{1}^{0}) = 130 \text{ GeV}, m(\tilde{\chi}_{1}^{0}) = 100 \text{ GeV} \\\Delta m(\tilde{\chi}_{2}^{0},\tilde{\chi}_{1}^{0}) = 130 \text{ GeV}, m(\tilde{\chi}_{1}^{0}) = 0 \text{ GeV}$	1908.03122 2103.08189
	squarks	~ ~ ~ .50	0-1 e, μ		E_T^{miss}	140	0.15-0.05	1.25 $m(\tilde{\chi}_2^0)=1$ GeV	2004.14060, 2012.03799
		$\tilde{\iota}_1 \tilde{\iota}_1, \tilde{\iota}_1 \rightarrow \iota \tilde{\chi}_1^0$ $\tilde{\iota}_1 \tilde{\iota}_1, \tilde{\iota}_1 \rightarrow W b \tilde{\chi}_1^0$	1 e,μ		E_T E_T^{miss}	140	Forbidden 1.0		2012.03799, 2401.13430
	gen. ect pr	$\tilde{i}_1 \tilde{i}_1, \tilde{i}_1 \rightarrow \tilde{\pi}_1 bv, \tilde{\pi}_1 \rightarrow \tau \tilde{G}$	1-2 τ	2 jets/1 b	E_T^{miss}	140	Forbidden	1.4 m(t)=800 GeV	2108.07665
	irec	$\tilde{t}_1 \tilde{t}_1, \tilde{t}_1 \rightarrow c \tilde{\chi}_1^0 / \tilde{c} \tilde{c}, \tilde{c} \rightarrow c \tilde{\chi}_1^0$	0 e, µ	2 c		36.1	0.85	$m(\tilde{\chi}_{\perp}^{0})=0$ GeV	1805.01649
•	d w		0 e,µ			140	0.55	$m(\tilde{t}_1,\tilde{c})-m(\tilde{\chi}_1^0)=5 \text{ GeV}$	2102.10874
		$ \begin{split} \tilde{t}_1 \tilde{t}_1, \tilde{t}_1 \rightarrow t \tilde{\chi}_2^0, \tilde{\chi}_2^0 \rightarrow Z/h \tilde{\chi}_1^0 \\ \tilde{t}_2 \tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + Z \end{split} $	1-2 e,μ 3 e,μ	1-4 b 1 b	E_T^{miss} E_T^{miss}	140 140	1 0.067- 52 Forbidden 0.86	1.18 $m(\tilde{\chi}_{2}^{0})=500 \text{ GeV}$ $m(\tilde{\chi}_{1}^{0})=360 \text{ GeV}, m(\tilde{r}_{1})-m(\tilde{\chi}_{1}^{0})=40 \text{ GeV}$	2006.05880 2006.05880
		${\tilde \chi}_1^\pm {\tilde \chi}_2^0$ via WZ	Multiple ℓ/jets ee, μμ	s ≥ 1 jet	E_T^{miss} E_T^{miss}	140 140	$\tilde{r}_{1}^{*}/\tilde{\chi}_{2}^{0}$ 0.205	$m(\tilde{\chi}_1^0)$ =0, wino-bino $m(\tilde{\chi}_1^0)$ =5 GeV, wino-bino	2106.01676, 2108.07586 1911.12606
		$\tilde{\chi}_1^{\pm} \tilde{\chi}_1^{\mp}$ via WW	2 e, µ		E_T^{miss}	140	0.42	$m(\tilde{\ell}_1^0)=0$, wino-bino	1908.08215
		$\tilde{\chi}_1^{\pm} \tilde{\chi}_2^0$ via Wh	Multiple ℓ/jets	6	E_T^{miss}		$\tilde{\chi}_1^{\pm}/\tilde{\chi}_2^0$ Forbidden 1.0		2004.10894, 2108.07586
		$\tilde{\chi}_1^{\pm} \tilde{\chi}_1^{\mp}$ via $\tilde{\ell}_L / \tilde{\nu}$	2 e, µ		E_T^{miss}	140	1.0	$m(\tilde{\ell}, \tilde{\nu})=0.5(m(\tilde{\chi}_1^{\pm})+m(\tilde{\chi}_1^{0}))$	1908.08215
	EW	$\tilde{\tau}\tilde{\tau}, \tilde{\tau} \rightarrow \tau \tilde{\chi}_1^0$	2τ 2e,μ	0 ioto	Emiss	140 140	[†] [[†] _R [†] _R] 0.35 0.5 0.7	$m(\tilde{x}_{1}^{0})=0$ $m(\tilde{x}_{1}^{0})=0$	2402.00603 1908.08215
	10	$\tilde{\ell}_{L,R} \tilde{\ell}_{L,R}, \tilde{\ell} \rightarrow \ell \tilde{\chi}_1^0$	ee, µµ	0 jets ≥ 1 jet	$\begin{array}{c} E_T^{\rm miss} \\ E_T^{\rm miss} \end{array}$	140	0.26	$m(\tilde{\ell})-m(\tilde{\ell}_1^0)=10 \text{ GeV}$	1911.12606
		$\tilde{H}\tilde{H}, \tilde{H} \rightarrow h\tilde{G}/Z\tilde{G}$	0 e, µ	$\geq 3 b$ 0 jets ≥ 2 large jets	E_T^{miss}	140	Ũ 0.94	$BR(\tilde{\chi}_{\perp}^{0} \rightarrow h\tilde{G})=1$	2401.14922
e			4 e,μ 0 e,μ	0 jets > 2 large jets	Emiss	140 140	0.55 0.45-0.93	$\begin{array}{c} BR(\hat{x}_{j}^{0} \rightarrow h\tilde{G}) = 1\\ BR(\hat{x}_{j}^{0} \rightarrow Z\tilde{G}) = 1\\ BR(\hat{x}_{1}^{0} \rightarrow Z\tilde{G}) = 1 \end{array}$	2103.11684 2108.07586
_			2 e,µ	≥ 2 jets	E_T^{miss}	140	й 0.77	$BR(\tilde{\chi}_1^0 \rightarrow Z\tilde{G})=BR(\tilde{\chi}_1^0 \rightarrow h\tilde{G})=0.5$	2204.13072
		$\operatorname{Direct} \tilde{\chi}_1^+ \tilde{\chi}_1^-$ prod., long-lived $\tilde{\chi}_1^\pm$	Disapp. trk	1 jet	$E_T^{\rm miss}$	140	0.66 0.21	Pure Wino Pure higgsino	2201.02472 2201.02472
	ved	Stable @ R-hadron	pixel dE/dx		Emiss	140	× 0.21	2.05	2205.06013
	-ong-lived particles	Metastable \tilde{g} R-hadron, $\tilde{g} \rightarrow qq \tilde{\chi}_1^0$	pixel dE/dx		$\begin{array}{c} E_T^{\rm miss} \\ E_T^{\rm miss} \end{array}$	140	ğ [τ(ĝ) =10 ns]	2.05 2.2 m($\tilde{\chi}_1^0$)=100 GeV	2205.06013
	ong	$\tilde{\ell}\ell, \tilde{\ell} \rightarrow \ell\tilde{G}$	Displ. lep		E_T^{miss}	140	. <i>ũ</i> 0.74	$r(\tilde{\ell}) = 0.1 \text{ ns}$	ATLAS-CONF-2024-011
	P		pixel dE/dx		E_T^{miss}	140	0.36 0.36	$ \begin{aligned} \tau(\tilde{\ell}) &= 0.1 \text{ ns} \\ \tau(\tilde{\ell}) &= 10 \text{ ns} \end{aligned} $	ATLAS-CONF-2024-011 2205.06013
		$\tilde{\chi}_1^{\pm} \tilde{\chi}_1^{\mp} / \tilde{\chi}_1^0, \tilde{\chi}_1^{\pm} \rightarrow Z \ell \rightarrow \ell \ell \ell$	3 e,µ	0.1.1	rumics	140	$\Gamma_{1}^{\mp}/\tilde{\lambda}_{1}^{0}$ [BR($Z\tau$)=1, BR(Ze)=1] 0.625 1.0		2011.10543
		$\tilde{\chi}_1^{\pm} \tilde{\chi}_1^{\mp} / \tilde{\chi}_2^0 \rightarrow WW / Z\ell\ell\ell\ell\nu\nu\nu$ $\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq \tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow qqq$	4 e,µ	0 jets ≥8 jets	E_T^{miss}	140 140	$\sum_{i=1}^{n+1} X_{2}^{0} = [\lambda_{i33} \neq 0, \lambda_{12k} \neq 0]$ 0.95 $\sum_{i=1}^{n} (m \bar{X}_{i}^{0}) = 50 \text{ GeV}, 1250 \text{ GeV}]$	1.55 m($\tilde{\chi}_1^0$)=200 GeV 1.6 2.34 Large χ_{11}''	2103.11684 2401.16333
	~	$gg, g \rightarrow qq\chi_1, \chi_1 \rightarrow qqq$ $\tilde{t}\tilde{t}, \tilde{t} \rightarrow t\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow tbs$		≥o jets Multiple		36.1	[m(x ₁)=50 GeV, 1250 GeV] [A ₂₃₃ =2e-4, 1e-2] 0.55 1.0		ATLAS-CONF-2018-003
	RPV	$\tilde{t}, \tilde{t} \rightarrow b\tilde{\chi}_{1}^{\pm}, \tilde{\chi}_{1}^{\pm} \rightarrow bbs$		$\geq 4b$		140	Forbidden 0.95	m($\tilde{\chi}_1^{\pm}$)=500 GeV	2010.01015
	Œ	$\tilde{t}_1 \tilde{t}_1, \tilde{t}_1 \rightarrow bs$		2 jets + 2 b		36.7	[1 [qq, bs] 0.42 0.61		1710.07171
		$\tilde{t}_1 \tilde{t}_1, \tilde{t}_1 \rightarrow q\ell$	2 e, µ	2 b		140	1 [1e-10< 1/32 <1e-8, 3e-10< 1/32 <3e-9] 1.0	0.4-1.85 BR(i ₁ → be/bµ)>20%	2406.18367
		$\tilde{\chi}_1^{\pm}/\tilde{\chi}_2^0/\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow tbs, \tilde{\chi}_1^{\pm} \rightarrow bbs$	1 μ 1-2 e, μ	DV ≥6 jets		136 140	$r_1 = [1e-10 < \lambda'_{23k} < 1e-8, 3e-10 < \lambda'_{23k} < 3e-9]$ 1.0 $r_2^0 = 0.2e-0.32$	1.6 $BR(\hat{i}_1 \rightarrow q\mu)=100\%, \cos\theta_i=1$ Pure higgsino	2003.11956 2106.09609
-		1,2							
									2

SUSY particle search

- Mass reach: near-maximal
- Displays the range in model space of search sensitivity.

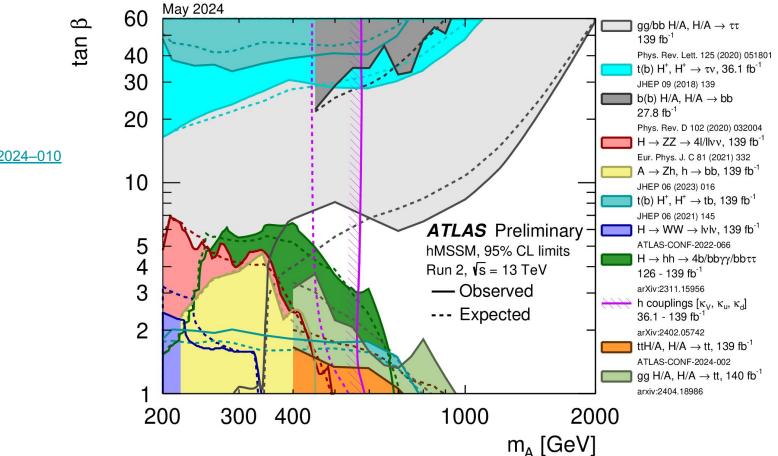
Note: some analyses have additional assumptions about intermediate states \rightarrow indicated by darker bands.

Long-lived sleptons: selectrons, smuons and staus with 0.3 ns lifetime are excluded, up to masses 740, 840 and 380 GeV, respectively

ATL-PHYS-PUB-2024-014

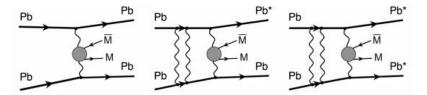
*Only a selection of the available mass limits on new states phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

Searches for SUSY Higgses



ATLAS-CONF-2024-010

Search for monopoles

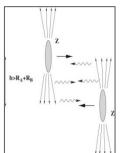


MoEDAL expt within CMS beam pipe: excluded magnetic monopoles with **mass < 80 GeV**

ATLAS ZDC: excluded mass < 120 GeV

2408.11035

2402.15628



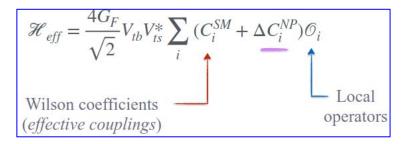
- Heavy ion collisions at the LHC generate magnetic fields stronger than in neutron stars.
- Such intense fields could lead to the spontaneous creation of composite magnetic monopoles.
- Ultra peripheral collisions (UPC) in lead-lead provide quasi-real photons→ suitable to search magnetic monopole pair production.
- Soft photons emitted by one lead nucleus can excite the other, typically through the giant dipole resonance,
 → induce the emission of one or more neutrons, each of which carries, on average, the full per-nucleon beam energy.
 - \Rightarrow three distinct EM breakup topologies:

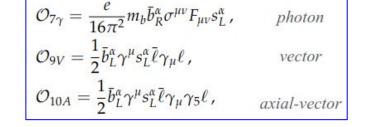
0n0n (no neutron emission) 0nXn (at least one neutron emitted by one nucleus) XnXn (at least one neutron is emitted by each nucleus) \rightarrow analysis uses only this data

- Use data at $\sqrt{s_{NN}}$ =5.36 TeV collected in 2023, with ZDC detector based software trigger
- Look at high pixel activity without associated reconstructed tracks, and low mass, low ET event

Amplitude analysis of $B^0 \rightarrow K^{*0}\mu^+\mu^-$ decays at LHCb

- Last few years: several anomalies in various measurements in heavy flavour sector (quark transitions observed in B hadron decays).
- Main issue: lower rate for $b \rightarrow s \mu^+ \mu^-$ in decays of B_s , B^0 , B^+
- Interpretation not straightforward due to hadronic uncertainties in SM predictions (form-factors, decay constant etc.).
- Also non-perturbative effects. Including long-distance charm loops.



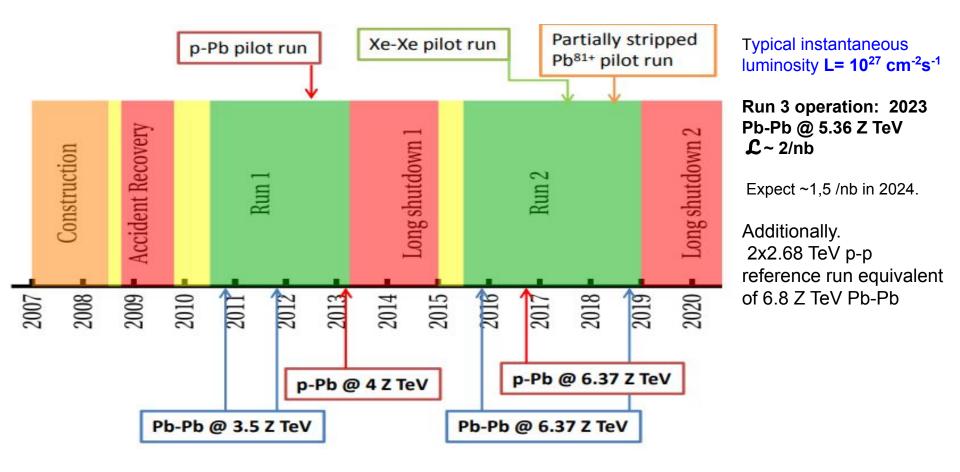


- First q²-unbinned (model-dependent fit) amplitude analysis of $B^0 \rightarrow K^{*0}\mu^+\mu^-$
- Estimation of non-local hadronic contributions (incl.charm loop) from data (with certain assumptions).
- Result consistent with anomalies observed in $b \rightarrow s \mu^+ \mu^-$ studies: 1.8 σ in C₉ and 1.4 σ global deviation in data from SM.

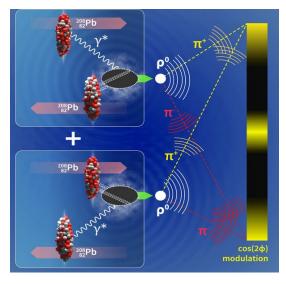
PRL 132 (2024) 13, 131801

 W^{-}

Heavy lon collisions at the LHC



Double-slit experiment by ALICE



Demonstration of the wave nature of propagating particles at the femtometre scale

Study vector meson production in ultra-peripheral collisions

 ρ^{0} produced within or close to one of the two well-separated nuclei.

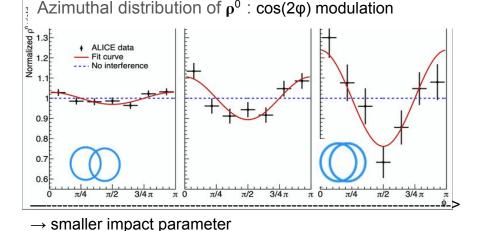
Decay rapidly: $\rho^0 \rightarrow \pi + \pi$ -, lifetime ~4.4* 10⁻²⁴ s

Symmetric system: cannot determine which of the nuclei emits the photon and which emits the two gluons => interference pattern akin to that of a double-slit interferometer.

different values of the impact parameter \rightarrow

 ϕ : angle between the two vectors formed by the sum and the difference of the transverse momenta of the pions.

The measured anisotropy corresponds to the amplitude of the $\cos(2\phi)$ modulation



2405.14525

Pb, p

Pb, p

o mesor

High Luminosity avatar of LHC: HL-LHC

- $\sqrt{s} = 14 \text{ TeV}$, ultimate L = **7.5X10³⁴ cm⁻²s⁻¹** (PU ~200)
- Expected total data vol. $\mathcal{L} \sim 4000 \text{ fb}^{-1}$
- **20 times more data** compared to current volume, in about **20 years time.**
- to be collected with **new avatars of the ATLAS and CMS** detectors.
- India playing significant role in the upgrade of several subsystems of CMS detector.

With more time achieve:

- More powerful analysis techniques
- More accurate theoretical tools
- Other "technological" breakthroughs (computing, AI, ...)
- New ideas!

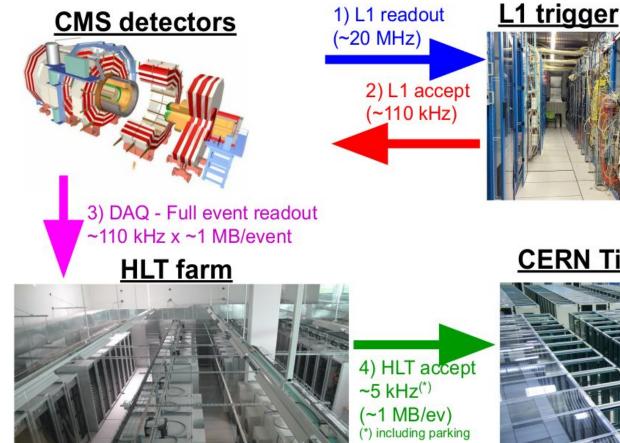
With more data study :

- Rare processes
- (Multi-) differential measurements
- Explore corners of phase space inaccessible with limited data volume



HL-LHC operation: 2029 -2040s

CMS experiment: Trigger Overview



Level-1 Trigger (L1T)

implemented in custom-designed electronics: high-end FPGAs, system-on-chip, high throughput data links.

Uses next generation control infrastructure based on Advanced Telecommunication Architecture (ATCA)

CERN Tier-0



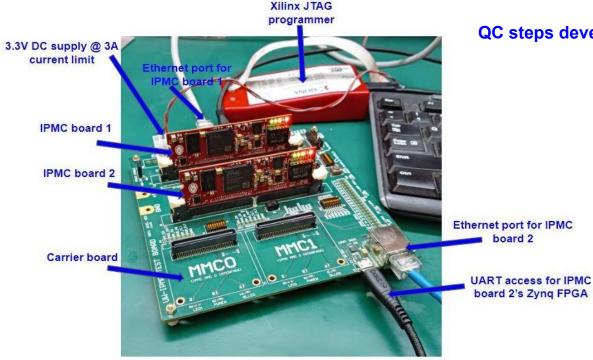


2

5) offline reconstruction

Indian contribution to hardware for L1 trigger upgrade

- Indian responsibility includes: hardware delivery + FPGA firmware development
- Mezzanine trigger boards: IPMC + ESM
- Fabrication in the Indian industry, quality control at TIFR.



QC steps developed, tools designed in TIFR

Hardware checks:

- Visual inspection of the edge / dimm connector under high resolution lens.
- LEDs powered on the board.

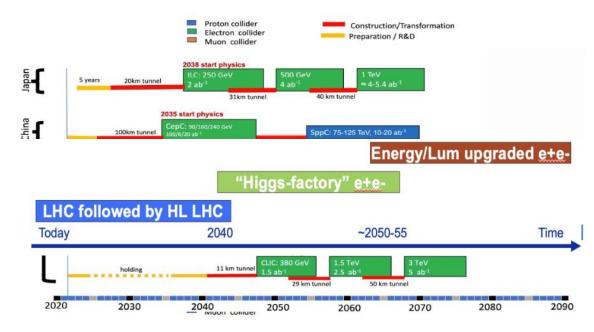
Key performance checks:

- Flash the bitstream through the Xilinx SDK tool.
- Boot image load into the flash using FTP.
- Run python software suite to check pin connectivity.

Next generation colliders on the plate

Data \Rightarrow significant gap between electroweak (EW) scale and the scale of New Physics (NP) \rightarrow use precision Higgs measurements as a tool to probe NP indirectly.

1% uncertainty in Higgs properties \Rightarrow 1 TeV scale of NP causing such a deviation \Rightarrow probe 10 TeV region \rightarrow go for exploratory hadron collider: 100 TeV FCC-hh!



In foreseeable future, mass-produce Higgs bosons in clean environment.

e+e- collider in about 15 -20 years time?

But Higgs factory is limited by lumi: can't probe rare H decays. Branching fractions vary over many orders of magnitude unlike for Z.

Summary

Presented glimpses on a very limited selection from a plethora of very interesting analyses carried out in recent times.

Run 3 data crucial to resolve some of the current disagreements between measurements and predictions; eg., cross sections for ttW, tWZ, tttt etc..

No direct evidence for physics beyond SM as yet.

However direct and indirect searches will continue.

More data allow us to look carefully in difficult corners of phasespace.

Need to study processes which are suppressed / forbidden in SM since New Physics interactions potentially enhance the rate providing access to higher mass scales in terms of virtual contribution.

New Physics can also effectively modify the couplings in various types of interactions

⇒ effective field theory (EFT) interpretation describes possible pattern of deviations introduced by new physics & also constrain the deviations.--> *did not discuss at all due to time constraint.*

The diversity of analyses indicates that the mining of interesting physics at the LHC will continue for next several decades.

LHC will always remain exciting and interesting!

Stay tuned!

2405.20040

Backup

EW results in two-photon collisions

RPP 87 (2024) 107801

OPAL

CMS

CMS

Photons can be simultaneously emitted by charged particles during p-p/ p-Pb/ Pb-Pb collisions

Cross section ~ $Z^4 \Rightarrow$ large rate of diphoton production in heavy ion runs.

At very high masses of $m_{\mu\nu}$, diffracted protons can be tagged by the forward detectors \Rightarrow study diffractive production of WW / $\tau\tau$

Utilize excellent tracking capability of experiments

CMS: ~ 30% of the 1 mm window around the beamspot not contain any pileup track.

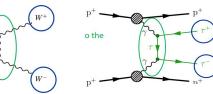
First observation of $\gamma\gamma \rightarrow \tau\tau$ in pp collisions by CMS

 $\sigma_{fid}(obs) = 11.2^{+3.1}_{-2.4} (syst)^{+2.2}_{-2.1} (stat) \text{ fb}, \text{ Significance 5.3 s.d. (6.5 exp.)}$

Constraints on the anomalous electromagnetic moments of T: a_ = 0.0009^{+ 0.0032} -0.0031

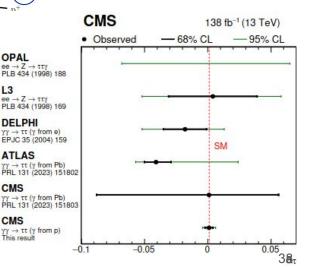
Dirac a₁: 0.0 Schwinger (SM) $a_{-} = 0.00116(9)$

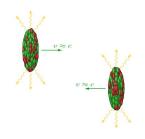
Dipole moment: $-1.7 < d_{\downarrow} < 1.7 \times 10^{-17}$ e cm.



2 back-to-back objects

- No hadronic activity close to the di-W/ τ vertex
- Ntracks = 0, pT > 0.5 GeV, $|\eta| < 2$





Exclusive production of high mass diphoton at 13 TeV

Also called light-by-light (LbyL) scattering .

Protons tagged in TOTEM precision proton spectrometer

Anomalous 4γ interaction in dim-8 EFT:

$$\mathcal{L}_{4\gamma} = \zeta_1 F_{\mu\nu} F^{\mu\nu} F_{\rho\sigma} F^{\rho\sigma} + \zeta_2 F_{\mu\nu} F^{\nu\rho} F_{\rho\lambda} F^{\lambda\mu},$$

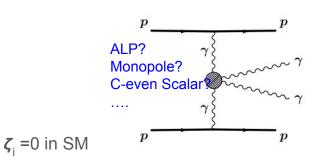
Study differential cross section:

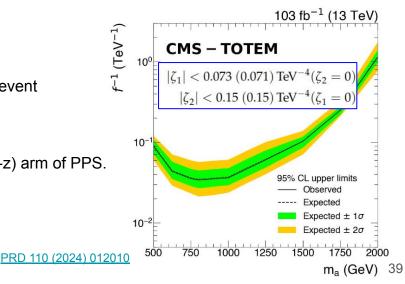
$$\frac{\mathrm{d}\sigma}{\mathrm{d}\Omega} = \frac{1}{16\pi^2 s} \left(s^2 + t^2 + st\right)^2 \left[48\zeta_1^2 + 40\zeta_1\zeta_2 + 11\zeta_2^2\right]$$

Only 1 event found in relevant phase-space, expected background: 1.1 event

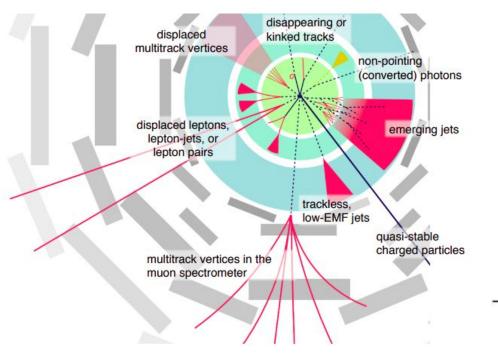
CMS $\sigma(pp \rightarrow p\gamma\gamma p) < 0.61 \text{ fb}$ for $p_T^{\gamma} > 100 \text{ GeV}$, $|\eta^{\gamma}| < 2.5$, $m_{\gamma\gamma} > 350 \text{ GeV}$ fractional proton energy loss of 0.035 < $\xi p < 0.150(0.180)$ for the +z (-z) arm of PPS.

Limits on axion-like particle (ALP) production in s-channel $\gamma\gamma \rightarrow a \rightarrow \gamma\gamma$ Coupling: $f^{-1} \ge 0.03$ to 1 TeV⁻¹ for m_a = 500–2000 GeV





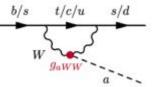
Search for long-lived particles



Possible candidates to explain neutrino mass,dark matter candidate

Unconventional signatures: require special reconstruction method, analysis strategy to identify displaced vertices.

Low mass neutral long-lived particles, including axion-like particles. Heavy neutral leptons Massive charged lon-lived particles Emerging jets



Signature: $a \rightarrow \gamma \gamma$ high energy deposit in electromagnetic calorimeter without any tracking Background: neutrino interaction with detector material

FASER experiment situated 500m from ATLAS collision point, aligned with beam collision axis.--> sensitive to long-lived axion-like-particles (ALPs) produced with (TeV) boost along the beam line and decaying inside the detector into photon pairs

Phenomenological MSSM interpretation

13 TeV data L = 138 /fb -> comprehensive analysis.

Generic realization of the MSSM with Lagrangian parameters defined at the supersymmetry (SUSY) scale **O**(1 TeV).

 \rightarrow captures most of the observable features of the general R-parity conserving weak scale MSSM

 \rightarrow allows more general conclusions to be drawn about SUSY compared with simplified models.

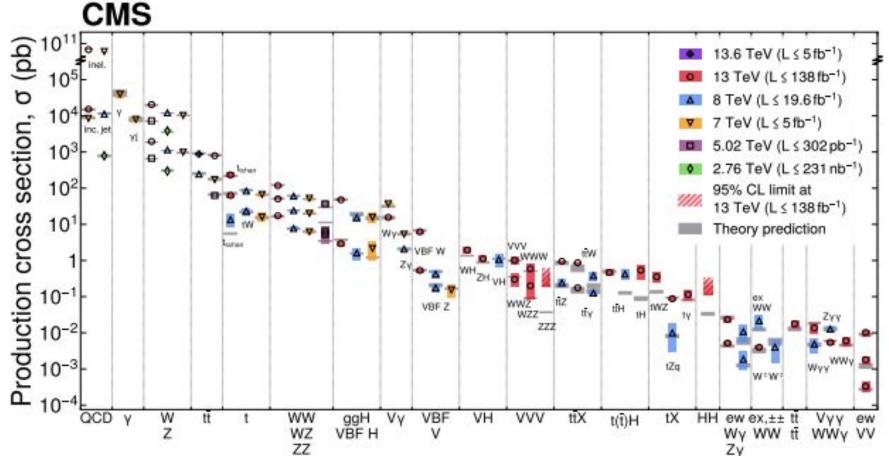
A global Bayesian analysis incorporates data from CMS as well as pre-CMS measuremnetsand indirect probes, estimating the marginalized posterior probability densities of model parameters, masses →observables based on the CMS results.

The CMS data highly suppress the phase space with colored superpartner masses below 1 TeV, considerably constrain natural SUSY and the electroweak sector, and weakly constrain SUSY dark matter.

- Significant phase space remains consistent with experimental data even at low LSP mass
- The lightest chargino, second-lightest neutralino, gluino, and top squark are heavily disfavored for masses less than around 200, 200, 700, and 1100 GeV, respectively.
- Considerable MSSM phase space capable of solving the small hierarchy problem or explaining the known DM relic density remain non-excluded by the CMS searches.
- only a very small number of models that are consistent with low-fine tuning and the relic density remain viable. Most such models correspond to a roughly pure Higgsino-like dark matter candidate.

Cross sections: from millibarn to femtobarn

2405.18661 Submitted to Phys.Rep.



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Physics model for W production and decay

